
LANDAUER'S PRINCIPLE FOR GAUSSIAN QUANTUM SYSTEMS

Gabriel T. Landi

Instituto de Física da Universidade de São Paulo

Caxambu d'Aju

May 29th, 2019



IFUSP

www.fmt.if.usp.br/~gtlandi

summary



- I. Landauer's principle.
- II. Entropy production.
- III. Quantum phase space.
- IV. Landauer for Gaussian states.

J. P. Santos, GTL and Mauro Paternostro, *Phys. Rev. Lett.*, **118**, 220601 (2017)

M. Brunelli, et. al., *Phys. Rev. Lett.*, **121**, 160604 (2018)

J. P. Santos, A. M. Timpanaro, M. Paternostro and GTL, *in preparation* (2019)

Landauer's principle: *information is physical*

- Fundamental **heat cost** for erasing information:

$$\Delta Q_E \geq -T \Delta S_S$$

ΔQ_E = heat dissipated in the reservoir.

ΔS_S = entropy change in the system/memory.

Mixed \rightarrow Pure: $\Delta S_S \leq 0 \rightarrow \Delta Q_E \geq 0$.



Landauer's principle: *information is physical*

- Fundamental **heat cost** for erasing information:

$$\Delta Q_E \geq -T \Delta S_S$$



ΔQ_E = heat dissipated in the reservoir.

ΔS_S = entropy change in the system/memory.

Mixed \rightarrow Pure: $\Delta S_S \leq 0 \rightarrow \Delta Q_E \geq 0$.

- $T = 0$: bound becomes uninformative.
- Would it be possible to have an erasure @ $T = 0$ with zero heat cost?

Landauer's principle: *information is physical*

- Fundamental **heat cost** for erasing information:

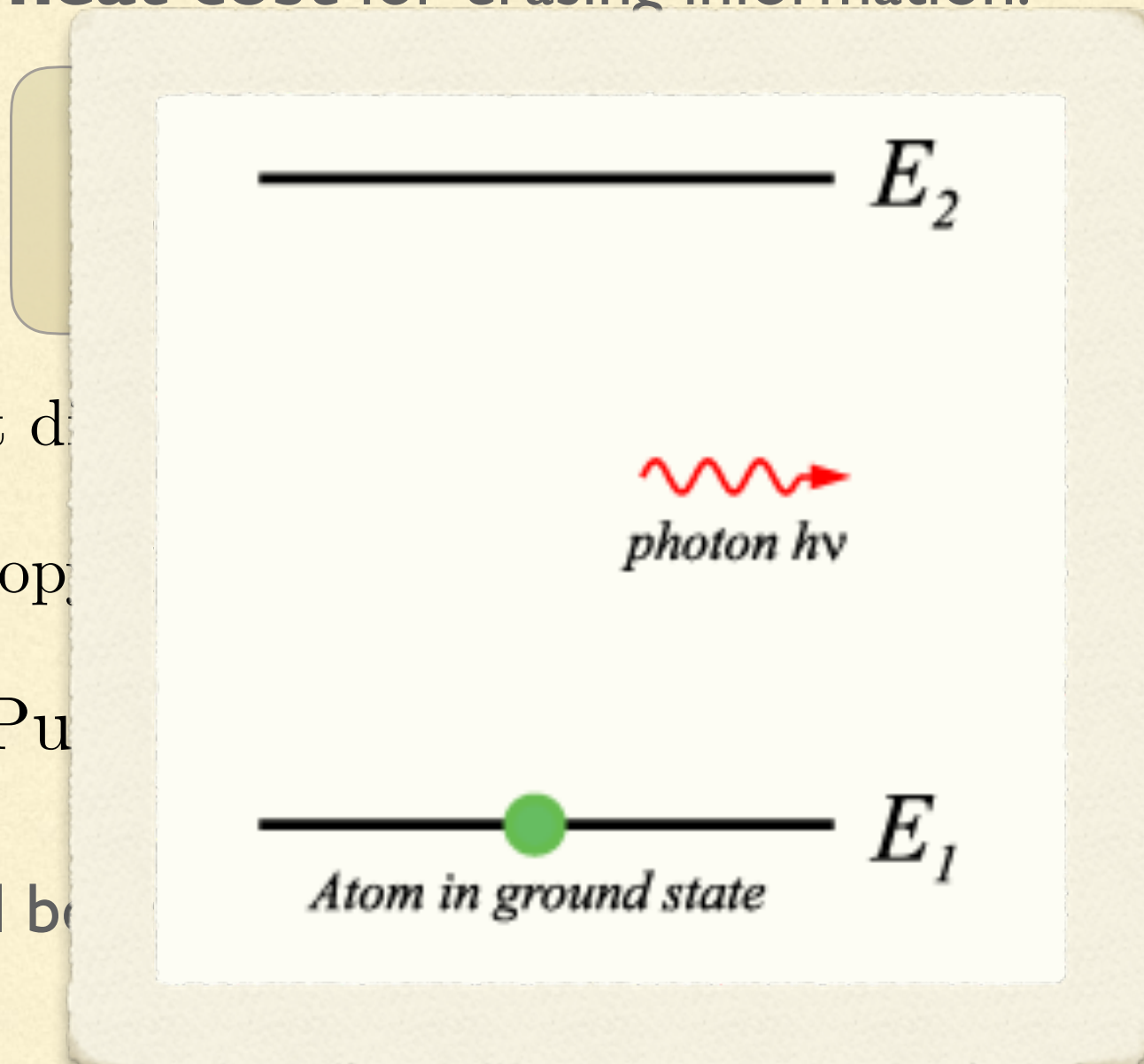
ΔQ_E = heat dissipated

ΔS_S = entropy change

Mixed \rightarrow Pure

- $T = 0$: bound by

- Would it be possible to have an erasure @ $T = 0$ with zero heat cost?



Open Quantum Dynamics

- System-environment interaction:

$$\rho'_{SE} = U(\rho_S \otimes \rho_E^{\text{th}})U^\dagger \quad \rightarrow \quad \rho'_S = \text{tr}_E \rho'_{SE}$$



Open Quantum Dynamics

- System-environment interaction:

$$\rho'_{SE} = U(\rho_S \otimes \rho_E^{\text{th}})U^\dagger \quad \rightarrow \quad \rho'_S = \text{tr}_E \rho'_{SE}$$

- Tracing over the environment is an irreversible process.
 - There is an associated entropy production:

$$\Sigma = \mathcal{I}'(S:E) + D(\rho'_E || \rho_E^{\text{th}})$$



Open Quantum Dynamics

- System-environment interaction:

$$\rho'_{SE} = U(\rho_S \otimes \rho_E^{\text{th}})U^\dagger \quad \rightarrow \quad \rho'_S = \text{tr}_E \rho'_{SE}$$



- Tracing over the environment is an irreversible process.
 - There is an associated entropy production:

$$\Sigma = \mathcal{I}'(S:E) + D(\rho'_E || \rho_E^{\text{th}})$$

$$\mathcal{I}'(S:E) = S(\rho'_S) + S(\rho'_E) - S(\rho'_{SE})$$

$$S(\rho) = -\text{tr} \rho \log \rho$$

$$D(\rho || \sigma) = \text{tr}(\rho \log \rho - \rho \log \sigma)$$

Open Quantum Dynamics

- System-environment interaction:

$$\rho'_{SE} = U(\rho_S \otimes \rho_E^{\text{th}})U^\dagger \quad \rightarrow \quad \rho'_S = \text{tr}_E \rho'_{SE}$$



- Tracing over the environment is an irreversible process.
 - There is an associated entropy production:

$$\Sigma = \mathcal{I}'(S:E) + D(\rho'_E || \rho_E^{\text{th}})$$

$$\Sigma \geq 0$$

2nd law

$$\mathcal{I}'(S:E) = S(\rho'_S) + S(\rho'_E) - S(\rho'_{SE})$$

$$S(\rho) = -\text{tr} \rho \log \rho$$

$$D(\rho || \sigma) = \text{tr}(\rho \log \rho - \rho \log \sigma)$$

Quantum Landauer's Principle

- It is possible to show that if the environment is thermal, then

$$\Sigma = \mathcal{I}'(S:E) + D(\rho'_E || \rho_E^{\text{th}}) = \beta \Delta Q_E - \Delta S_S \geq 0$$

- Landauer's principle is thus a *direct* consequence of the 2nd law.

Quantum Landauer's Principle

- It is possible to show that if the environment is thermal, then

$$\Sigma = \mathcal{I}'(S:E) + D(\rho'_E || \rho_E^{\text{th}}) = \beta \Delta Q_E - \Delta S_S \geq 0$$

- Landauer's principle is thus a *direct* consequence of the 2nd law.
- But this continues to be problematic @ $T = 0$.
 - Note that in this case the relative entropy diverges because the target state becomes pure.

Wigner entropy production

- Focus on continuous variable systems with Gaussian states (q-optics, BECs, phonons, trapped ions, &c.).
- Instead of using the von Neumann entropy, use the entropy of the *Wigner function*

$$S_W = - \int d^2\alpha \, W \log W$$

J. P. Santos, GTL and Mauro Paternostro, *Phys. Rev. Lett.*, **118**, 220601 (2017)
M. Brunelli, et. al., *Phys. Rev. Lett.*, **121**, 160604 (2018)

Wigner entropy production

- Focus on continuous variable systems with Gaussian states (q-optics, BECs, phonons, trapped ions, &c.).
- Instead of using the von Neumann entropy, use the entropy of the *Wigner function*

$$S_W = - \int d^2\alpha \, W \log W$$

- For Gaussian states, the Wigner function is always non-negative.
- Moreover, it is related to the Rényi-2 entropy:

$$S_2 = - \log \text{tr} \rho^2$$

J. P. Santos, GTL and Mauro Paternostro, *Phys. Rev. Lett.*, **118**, 220601 (2017)

M. Brunelli, et. al., *Phys. Rev. Lett.*, **121**, 160604 (2018)

Adesso, Girolami, Serafini, *PRL*, **109**, 190502 (2012)

Wigner-Landauer principle for CVs

- We assume that S and E are both bosonic and Gaussian.
 - Moreover, their unitary is Gaussian preserving.
- Define the Wigner entropy production:

$$\Sigma_W = \mathcal{I}'_W(S:E) + D_W(W'_E || W_E^{\text{th}}) \geq 0$$

- where all quantities are now defined in terms of the Wigner function. e.g.,

$$D_W(W_1 || W_2) = \int d^2\alpha \, W_1 \log W_1 / W_2$$

Wigner-Landauer principle for CVs

- We assume that S and E are both bosonic and Gaussian.
 - Moreover, their unitary is Gaussian preserving.
- Define the Wigner entropy production:

$$\Sigma_W = \mathcal{I}'_W(S:E) + D_W(W'_E || W_E^{\text{th}}) \geq 0$$

- where all quantities are now defined in terms of the Wigner function. e.g.,

$$D_W(W_1 || W_2) = \int d^2\alpha \, W_1 \log W_1 / W_2$$

- For Gaussian states, the Wigner entropy satisfies the strong subadditivity inequality. All entropic quantities are well behaved.
-

-
- When the environment is thermal, we can also write this as

$$\Sigma_W = \Delta\Phi_E - \Delta S_{W,S} \geq 0$$

-
- When the environment is thermal, we can also write this as

$$\Sigma_W = \Delta\Phi_E - \Delta S_{W,S} \geq 0 \qquad \Delta\Phi_E = \sum_{k=1}^{N_E} \frac{\langle b_k^\dagger b_k \rangle' - \langle b_k^\dagger b_k \rangle_{\text{th}}}{\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2}$$

-
- When the environment is thermal, we can also write this as

$$\Sigma_W = \Delta\Phi_E - \Delta S_{W,S} \geq 0$$
$$\Delta\Phi_E = \sum_{k=1}^{N_E} \frac{\langle b_k^\dagger b_k \rangle' - \langle b_k^\dagger b_k \rangle_{\text{th}}}{\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2}$$
$$\langle b_k^\dagger b_k \rangle_{\text{th}} = \frac{1}{e^{\Omega_k/T} - 1}$$

-
- When the environment is thermal, we can also write this as

$$\Sigma_W = \Delta\Phi_E - \Delta S_{W,S} \geq 0 \qquad \Delta\Phi_E = \sum_{k=1}^{N_E} \frac{\langle b_k^\dagger b_k \rangle' - \langle b_k^\dagger b_k \rangle_{\text{th}}}{\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2}$$
$$\langle b_k^\dagger b_k \rangle_{\text{th}} = \frac{1}{e^{\Omega_k/T} - 1}$$

- This is not yet very useful:

- What makes Landauer useful is the connection to **heat**:

$$\Delta Q_E = \sum_k \Omega_k (\langle b_k^\dagger b_k \rangle' - \langle b_k^\dagger b_k \rangle_{\text{th}})$$

- When the environment is thermal, we can also write this as

$$\Sigma_W = \Delta\Phi_E - \Delta S_{W,S} \geq 0 \qquad \Delta\Phi_E = \sum_{k=1}^{N_E} \frac{\langle b_k^\dagger b_k \rangle' - \langle b_k^\dagger b_k \rangle_{\text{th}}}{\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2}$$

$$\langle b_k^\dagger b_k \rangle_{\text{th}} = \frac{1}{e^{\Omega_k/T} - 1}$$

- This is not yet very useful:

- What makes Landauer useful is the connection to **heat**:

$$\Delta Q_E = \sum_k \Omega_k (\langle b_k^\dagger b_k \rangle' - \langle b_k^\dagger b_k \rangle_{\text{th}})$$

- Playing with a chain of inequalities leads to:

$$\Delta Q_E \geq \Gamma_{\min} \Delta S_{W,S}, \qquad \Gamma_{\min} = \min_k \Omega_k (\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2)$$

$$\Delta Q_E \geq \Gamma_{\min} \Delta S_{W,S}, \quad \Gamma_{\min} = \min_k \Omega_k (\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2)$$

$$\Delta Q_E \geq \Gamma_{\min} \Delta S_{W,S}, \quad \Gamma_{\min} = \min_k \Omega_k (\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2)$$

- For high temperatures ($T \gg \Omega_k, \quad \forall k$)

$$\Omega_k (\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2) \simeq T$$

- We then recover the usual Landauer principle.
-

$$\Delta Q_E \geq \Gamma_{\min} \Delta S_{W,S}, \quad \Gamma_{\min} = \min_k \Omega_k (\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2)$$

- For high temperatures ($T \gg \Omega_k, \quad \forall k$)

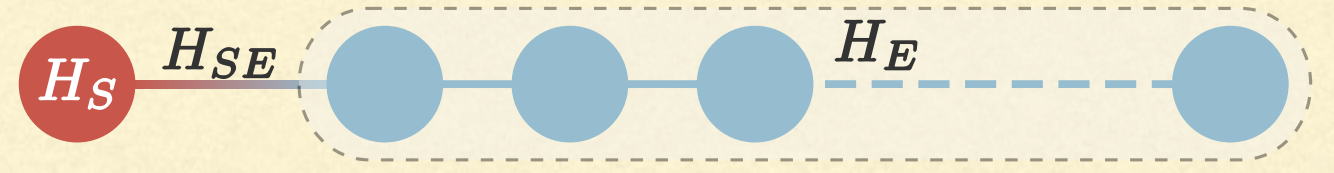
$$\Omega_k (\langle b_k^\dagger b_k \rangle_{\text{th}} + 1/2) \simeq T$$

- We then recover the usual Landauer principle.
- But now if $T = 0$

$$\Gamma_{\min} = \min_k \Omega_k / 2 = \text{infrared cutoff}$$

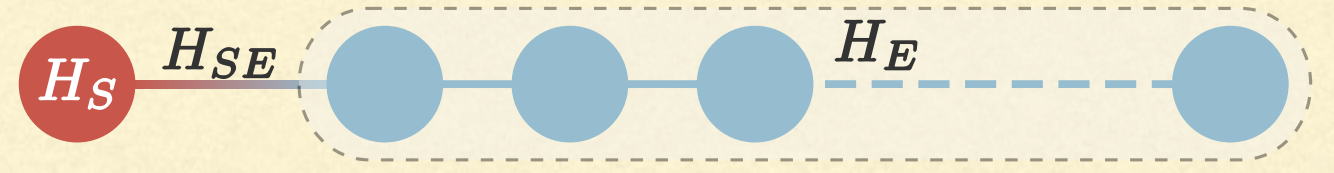
- The infrared cutoff establishes a fundamental lower bound at zero temperature
-

Example: linear chain



$$H_{SE} = \omega a^\dagger a + \sum_{i=1}^{N_E} \omega_0 c_i^\dagger c_i - g \sum_{i=1}^{N_E-1} (c_i^\dagger c_{i+1} + c_{i+1}^\dagger c_i) - \lambda(a^\dagger c_1 + a c_1^\dagger)$$

Example: linear chain



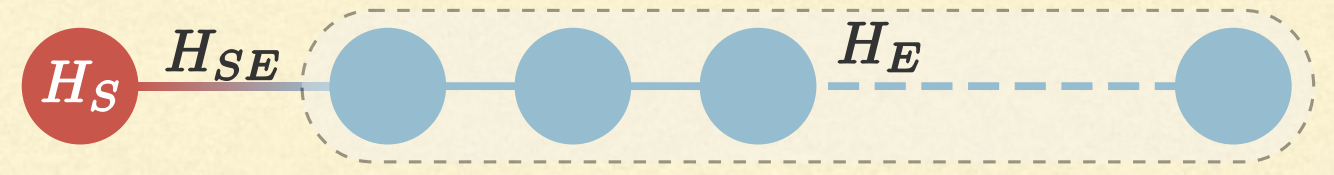
$$H_{SE} = \omega a^\dagger a + \sum_{i=1}^{N_E} \omega_0 c_i^\dagger c_i - g \sum_{i=1}^{N_E-1} (c_i^\dagger c_{i+1} + c_{i+1}^\dagger c_i) - \lambda(a^\dagger c_1 + a c_1^\dagger)$$

- Diagonalizing the bath part leads to normal modes:

$$b_k = \sqrt{\frac{2}{N_E + 1}} \sum_{i=1}^{N_E} \sin(ik) c_i$$

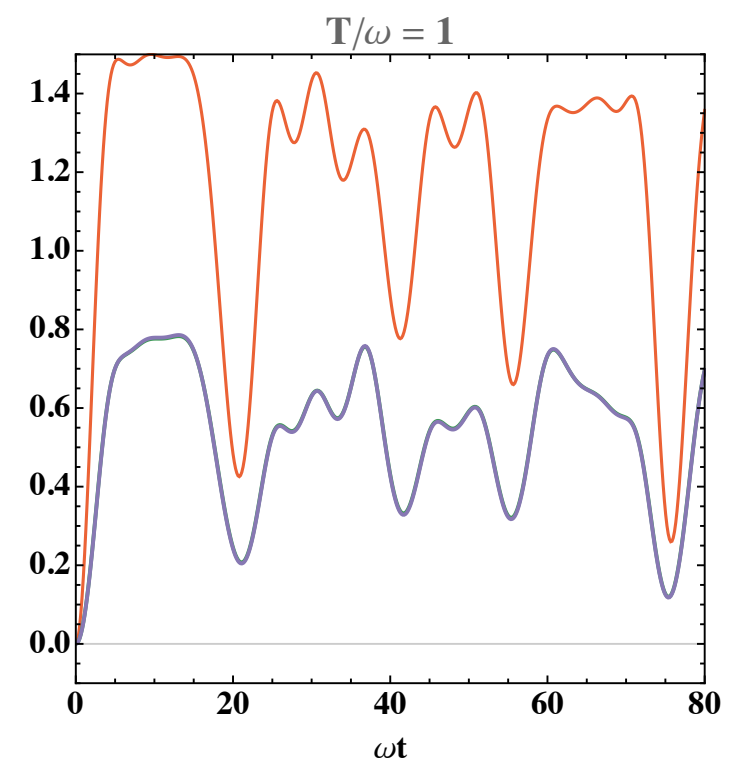
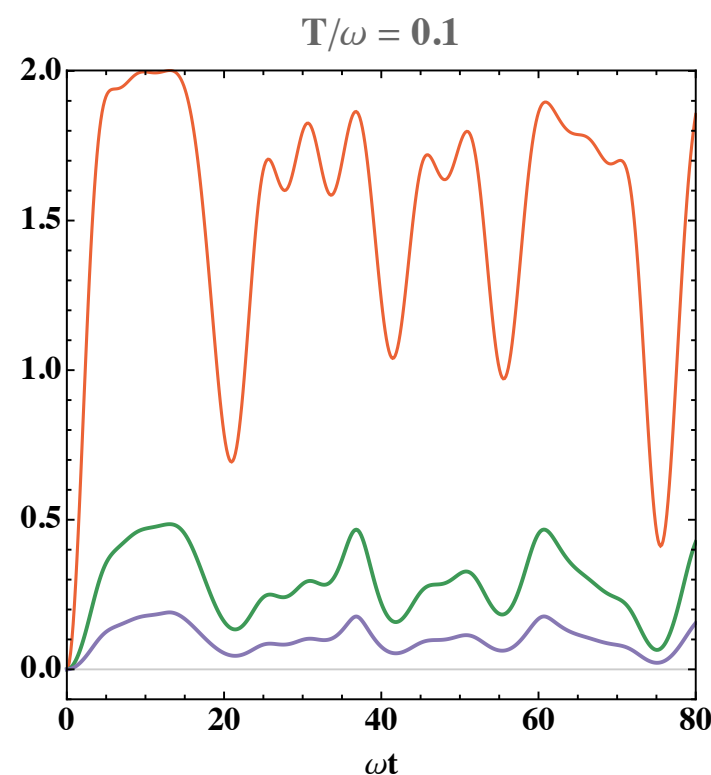
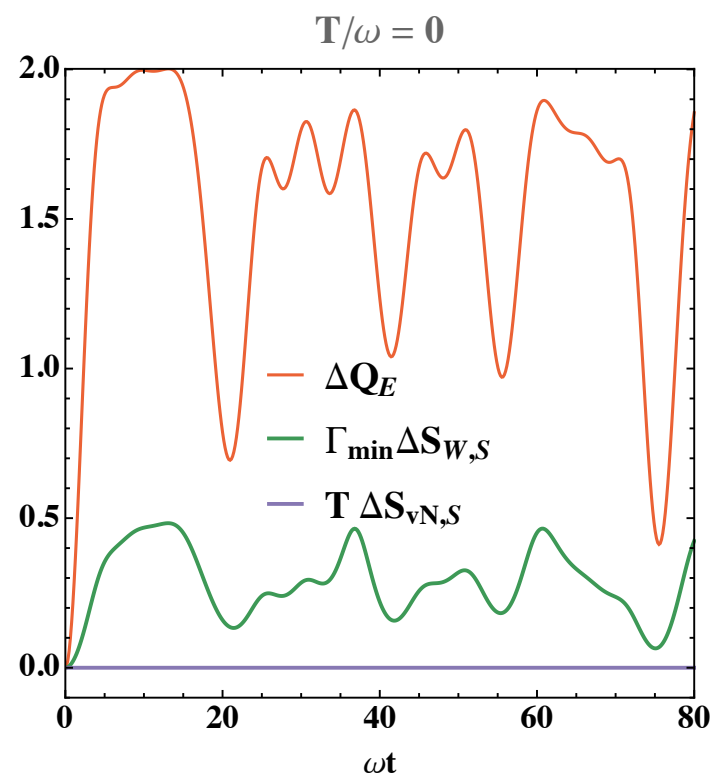
$$\Omega_k = \omega_0 - g \cos(k)$$

Example: linear chain

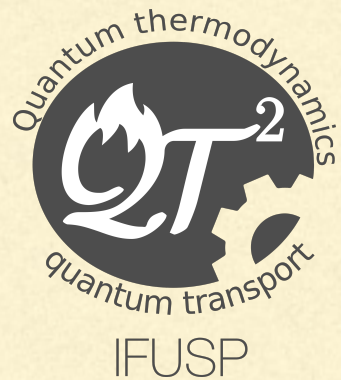


$$H_{SE} = \omega a^\dagger a + \sum_{i=1}^{N_E} \omega_0 c_i^\dagger c_i - g \sum_{i=1}^{N_E-1} (c_i^\dagger c_{i+1} + c_{i+1}^\dagger c_i) - \lambda(a^\dagger c_1 + a c_1^\dagger)$$

- Diagonalizing the bath part leads to normal modes:



Conclusions

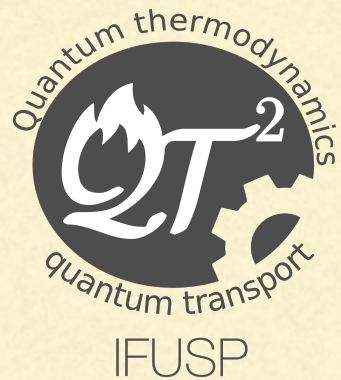


www.fmt.if.usp.br/~gtlandi

Acknowledgements:
IFUSP, FAPESP, CNPq

Conclusions

- Landauer's principle: *heat up to erase information*

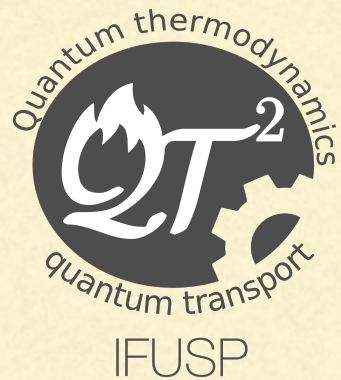


www.fmt.if.usp.br/~gtlandi

Acknowledgements:
IFUSP, FAPESP, CNPq

Conclusions

- Landauer's principle: *heat up to erase information*
- Landauer is uninformative when $T = 0$.

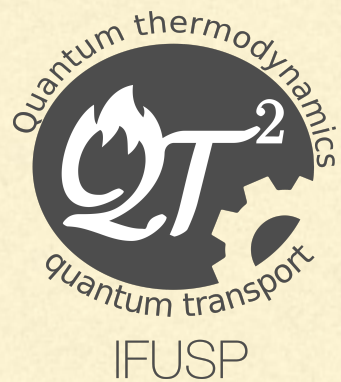


www.fmt.if.usp.br/~gtlandi

Acknowledgements:
IFUSP, FAPESP, CNPq

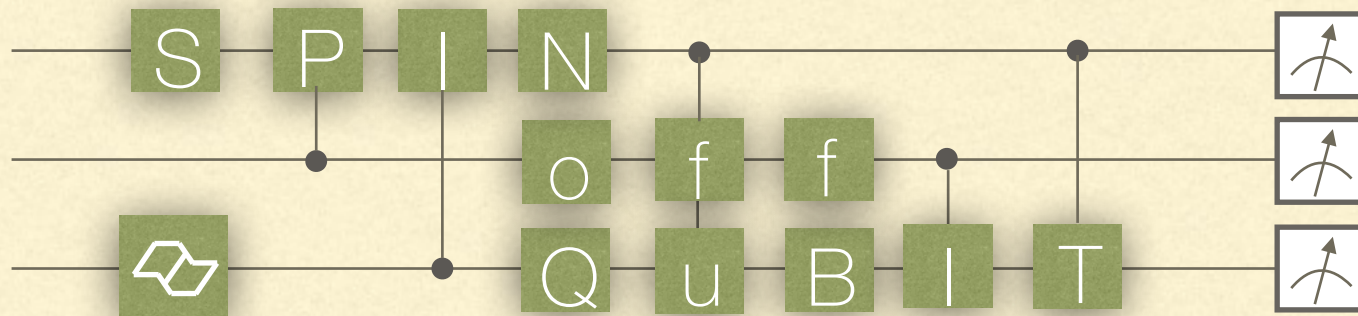
Conclusions

- Landauer's principle: *heat up to erase information*
- Landauer is uninformative when $T = 0$.
- Formulating the problem in terms of phase space entropies reveals a minimum heat cost.

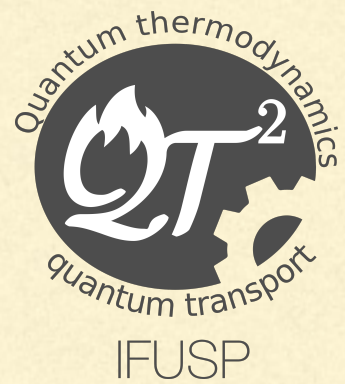


www.fmt.if.usp.br/~gtlandi

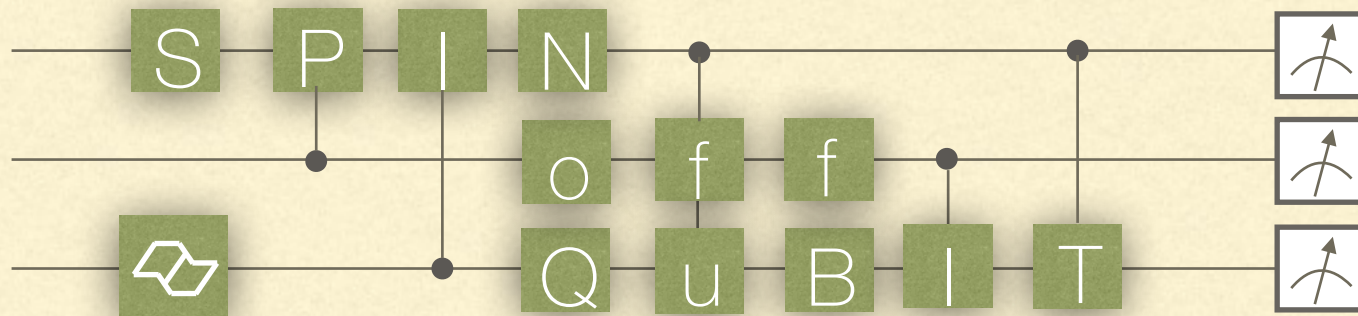
Acknowledgements:
IFUSP, FAPESP, CNPq



spinoffqubit.info



www.fmt.if.usp.br/~gtlandi



Thank you!

spinoffqubit.info



www.fmt.if.usp.br/~gtlandi