

Energy barriers between metastable states in first order quantum phase transitions

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S Wald, A Timpanaro, C Cormick and GT Landi



Queens, Feb. 15, 2018

Bose Hubbard Model (BHM)

Gersch, Knollman 1963

$$H_{\text{BH}} = \mathcal{T} + \sum_{i=1}^K \frac{U_s}{2} n_i(n_i - 1) - \mu \sum_{i=1}^K n_i$$

$\mathcal{T} = -J \sum_{\langle i,j \rangle} (b_i^\dagger b_j + b_j^\dagger b_i)$

nearest neighbor hopping

on site repulsion

chemical potential

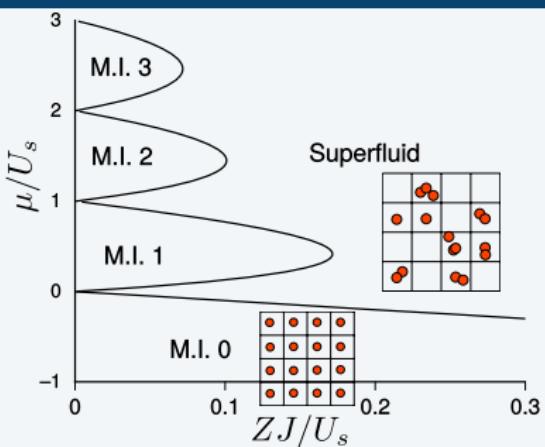
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mean-field phase diagram Sachdev

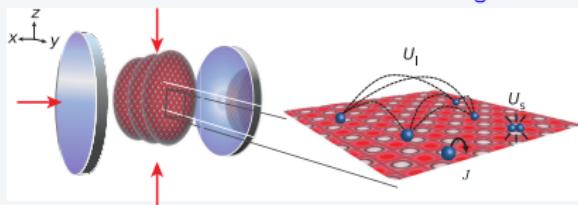


applications

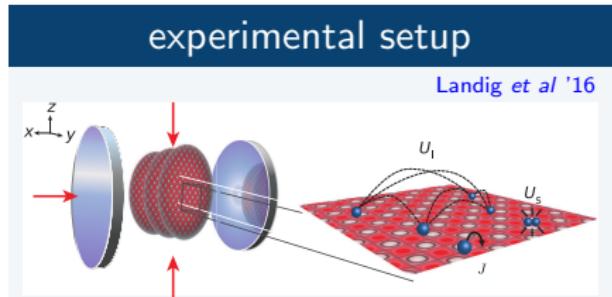
- solvable in MF (\leftrightarrow Fermi-HM)
- captures MI - SF transition
- ultra-cold atoms in optical lattices
- next step after BEC towards macroscopic quantum phenomena
- connect to many-body quantum properties

experimental setup

Landig et al '16



- BEC: $4.2(4) \cdot 10^4$ ^{87}Ru atoms
- 2D optical lattice:
- **ultrahigh-finesse optical cavity**
- ~ 2.8 atoms per lattice site



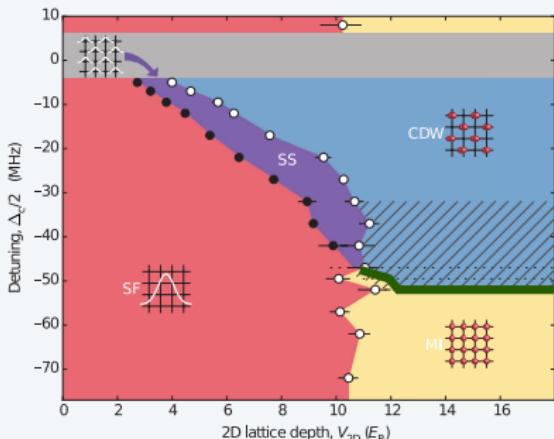
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theoretical description (canonical)

$$H = H_{\text{BH}} - \frac{U_\ell}{K} \left(\underbrace{\sum_{i \in e} n_i - \sum_{i \in o} n_i}_{\Theta} \right)^2$$

experimental results

Landig et al, [Nature'16](#)



Long range BHM

Landig et al Nature'16

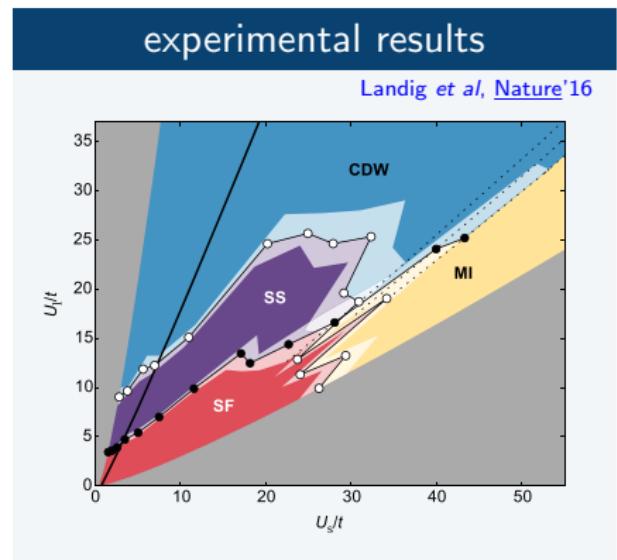
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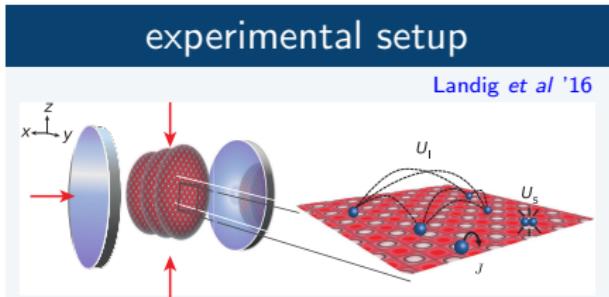
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Long range BHM

Landig et al Nature'16



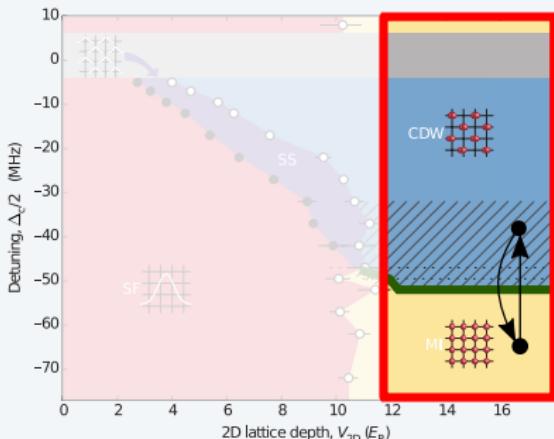
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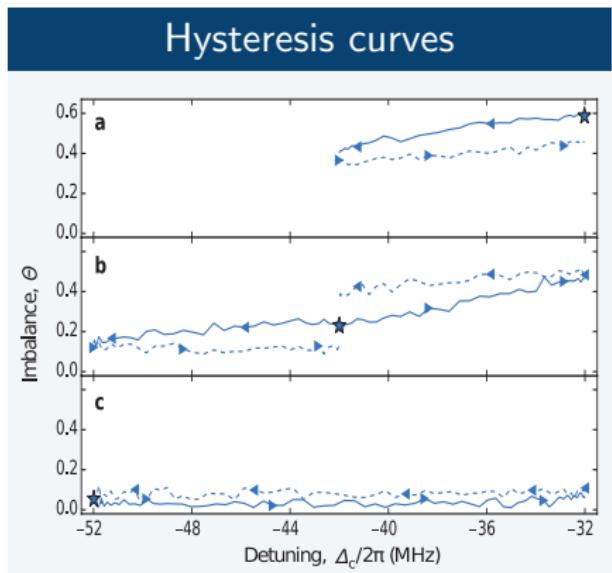
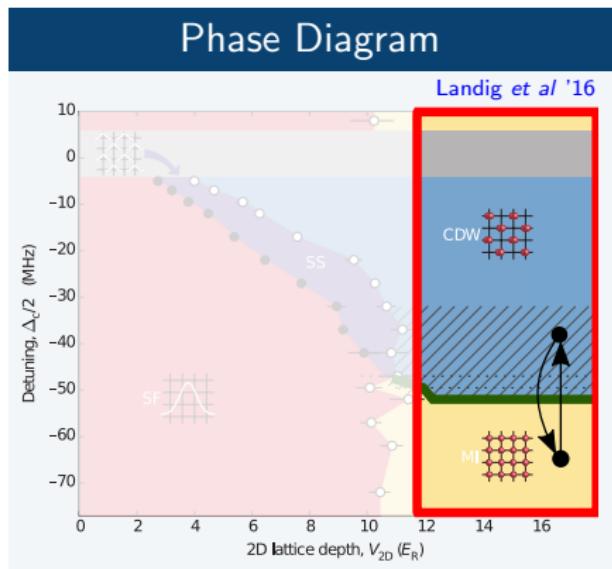
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experimental results

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Hysteresis



Overview

- 1 Exact Landau theory in the zero hopping limit
 - Energy landscape and hysteresis loops for $\rho = 1$
 - Different fillings: $\rho \neq 1$
- 2 Variational Ansatz: $J \neq 0$
 - Physical requirements for reduced Hilbert space
 - Construction of variational states
- 3 Numerical Analysis
 - Physical observables and phase diagram
- 4 Discrete WKB method
 - Energy barriers between meta-stable quantum phases
- 5 Conclusion
- 6 Outlook

The Zero Hopping Limit

Landau free energy

- CDW & MI: $J/U_s \ll 1$

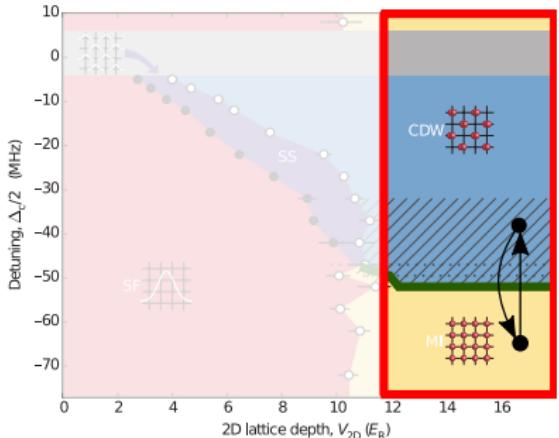
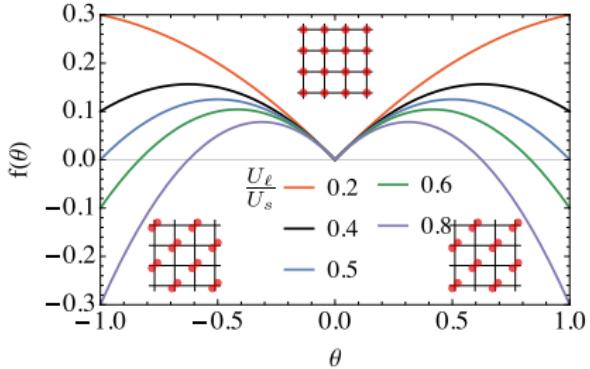
$$H_{J \rightarrow 0} = \sum_{i=1}^K \frac{U_s}{2} \hat{n}_i (\hat{n}_i - 1) - \frac{U_\ell}{K} \Theta^2$$

- Fock basis: H diagonal
- free energy:

$$f(\theta) = -U_\ell \theta^2 + \frac{\phi(\rho + \theta) + \phi(\rho - \theta)}{2}$$

$$\phi(\rho_x) = \frac{U_s}{K} \min_{N_x} \left\{ \sum_i n_i (n_i - 1) \right\}$$

- $\rho = 1 \Rightarrow f(\theta) = -U_\ell \theta^2 + \frac{U_s}{2} |\theta|$



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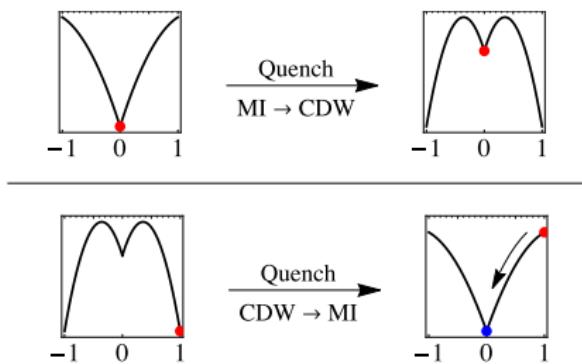
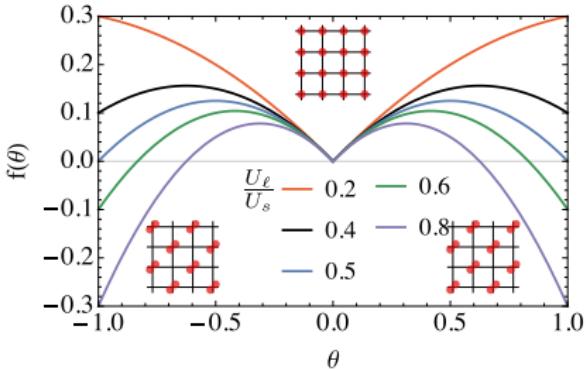
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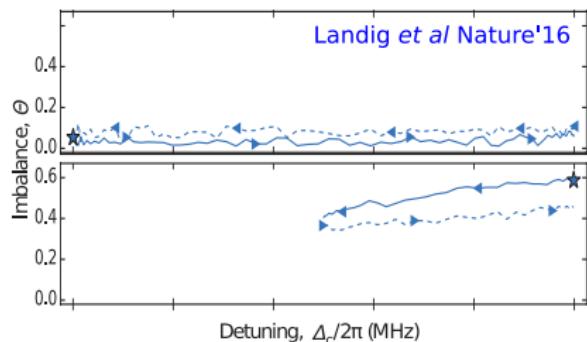
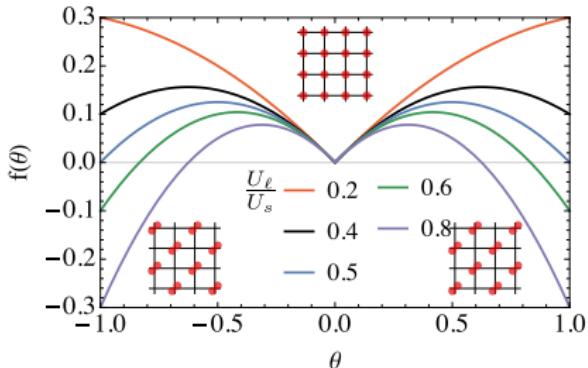
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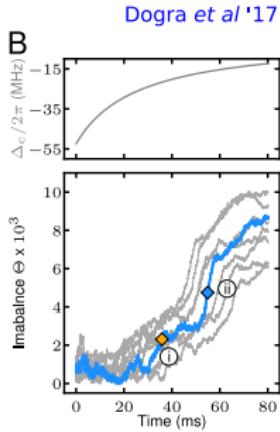
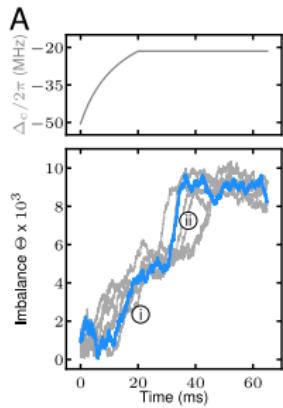
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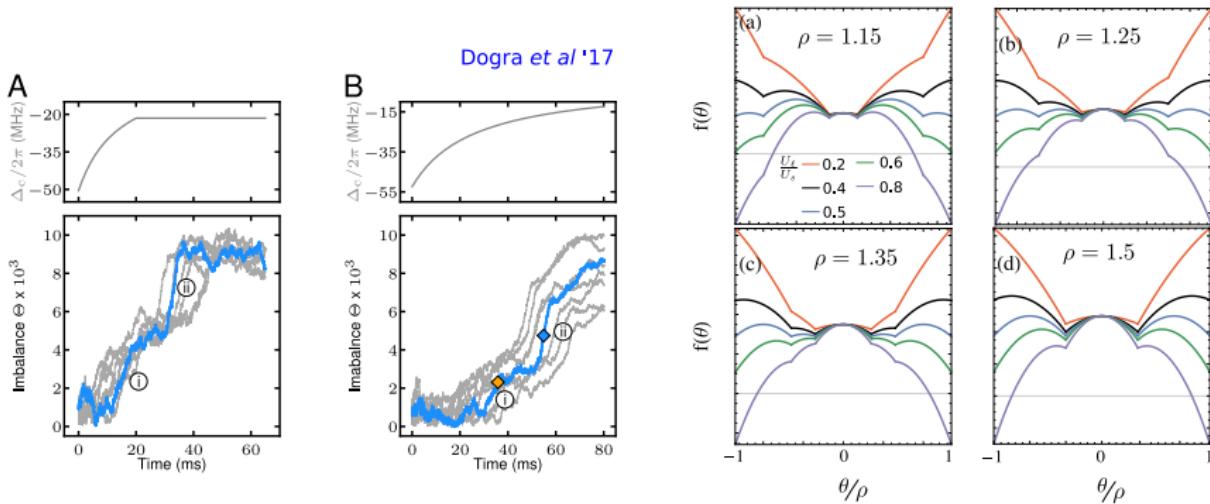
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Different fillings: $\rho \neq 1$



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- several meta-stable minima appear
- qualitative explanation for plateaus in experiments

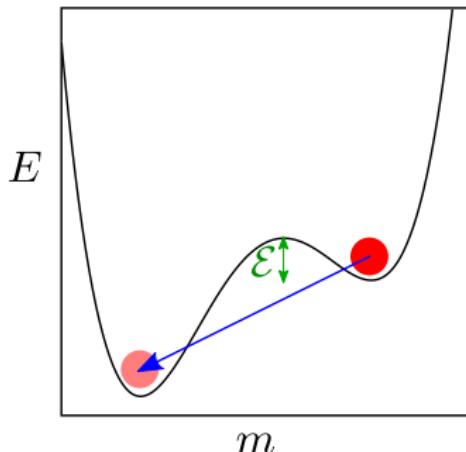
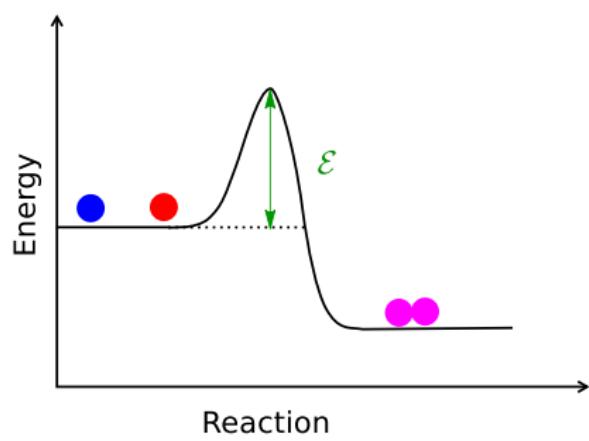
core message

energy landscape explains many experimentally relevant findings

Energy barriers and landscapes

Goal

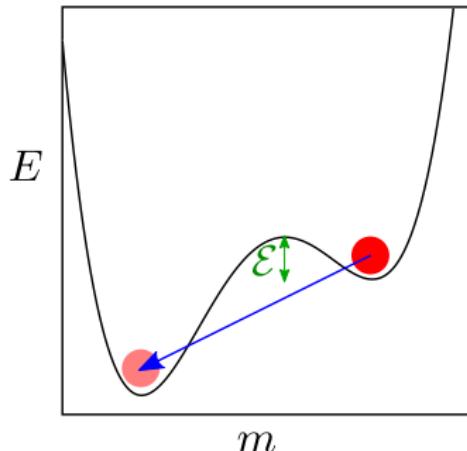
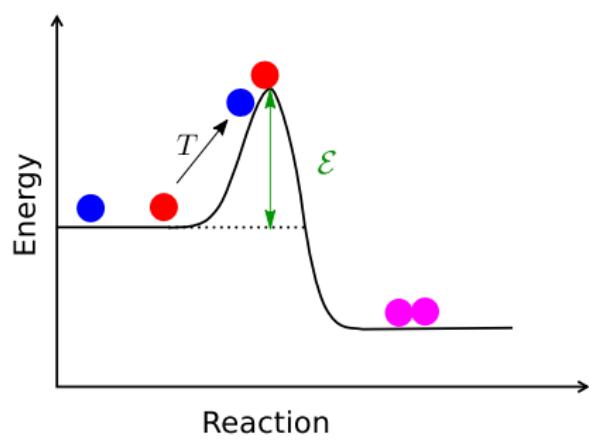
- describe idea of an energy barrier between the MI & CDW
- how to overcome barrier ?
- Arrhenius theory, Ginzburg-Landau theory: thermally assisted
- $T = 0$? quantum fluctuations ?



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Variational Ansatz: $J \neq 0$

Requests

- CDW & MI: $J/U_s \ll 1$
- focus on $\rho = 1$
- $|MI\rangle = |1\dots1; 1\dots1\rangle$
 $|CDW\rangle = |2\dots2; 0\dots0\rangle$
- choose intermediate states:
 - restrict to $n_i = 0, 1, 2 \ \forall i$
 - $\Theta|Q, \nu\rangle = Q|Q, \nu\rangle$ with
 $\Theta = (\sum_{i \in e} n_i - \sum_{i \in o} n_i)^2$
- $|1, 2, 1, 2; 0, 1, 0, 1\rangle; |1, 2, 1, 2; 0, 0, 0, \textcolor{red}{2}\rangle$
 - additional U_s without U_ℓ gain
 - distribute atoms in sublattices

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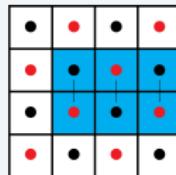
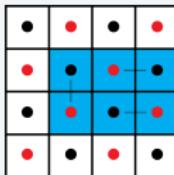
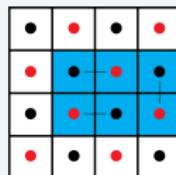
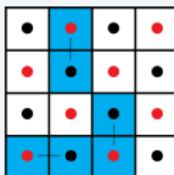
Choice of variational states

- $\mathcal{T} = -\sqrt{2}J(\mathcal{T}_e + \mathcal{T}_o)$, with
 - (i) $\mathcal{T}_o = \mathcal{T}_e^\dagger$
 - (ii) $[\Theta, \mathcal{T}_e] = 2\mathcal{T}_e$
 - (iii) $[\Theta, \mathcal{T}_o] = -2\mathcal{T}_o$
- $\mathcal{T}_{e/o}$: creation/annihilation operator of imbalance
- tight-binding:
$$\langle \psi | H | \psi \rangle = \sum_Q \epsilon_Q \psi_Q^* \psi_Q + \gamma_Q^+ \psi_{Q+2}^* \psi_Q + \gamma_Q^- \psi_{Q-2}^* \psi_Q$$
- nucleation of CDW 0 – 2 pairs
- $|Q\rangle = \frac{1}{\sqrt{A(Q)}} (\tilde{\mathcal{T}}_e)^{Q/2} |MI\rangle$

Variational Ansatz - Mean Field

Normalisation constants

- $\gamma_Q^+ := \langle Q + 2|T|Q \rangle \propto \sqrt{\frac{A(Q+2)}{A(Q)}}$
- normalisation constants?
- maps to matching problem



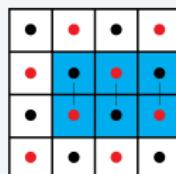
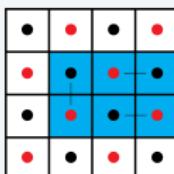
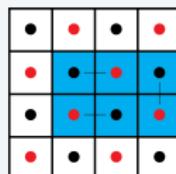
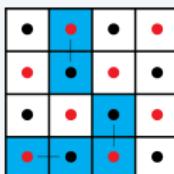
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- # P-hard $\leftrightarrow Z_{3D-Ising}$

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- **normalisation constants?**
- maps to matching problem



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- # P-hard $\leftrightarrow Z_{3D-Ising}$

Lattice deformation

- deform CDW generator
- **eliminate lattice structure**
- ∞ -range hopping $e \leftrightarrow o$
- emergence of **non-local** CDW

- for Q -states only !

- $A(Q) = \binom{K/2}{Q/2}^2$

$$\gamma_Q^+ = -\frac{\alpha}{4K} \begin{cases} (K-Q)(Q+2) & Q \geq 0 \\ (K-|Q|+2)|Q| & Q < 0 \end{cases}$$

- **no predictions** for lattice dependent properties (e.g. CDW-SF)

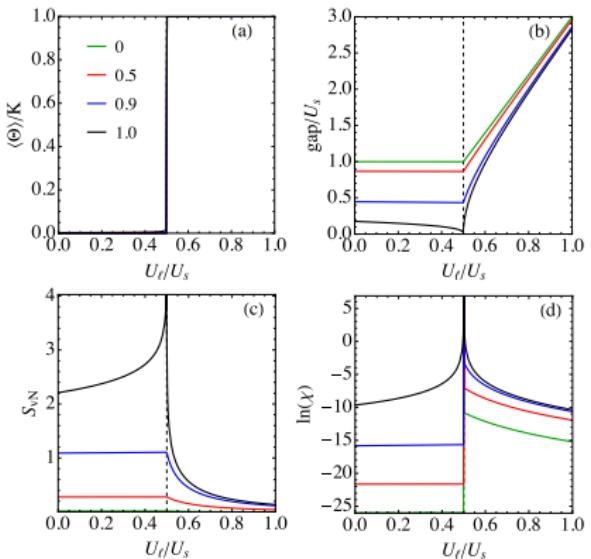
Numerical Analysis

Numerical details

- analyse GS properties
- $K = 2000$ lattice sites
- different hopping: $\alpha = 8\sqrt{2}J$

Physical quantities

- $\langle \Theta \rangle / K = \pm 1$: CDW or MI ?
- $\Delta E_{k \rightarrow 0}$: compressible phase ?
- $S_{vN} = - \text{tr } \rho_e \ln \rho_e$
- $\chi = -\partial_\delta^2 \ln |\langle \psi(U_\ell) | \psi(U_\ell + \delta) \rangle| \Big|_{\delta=0}$
pinpoints all transitions for $K = 10$!



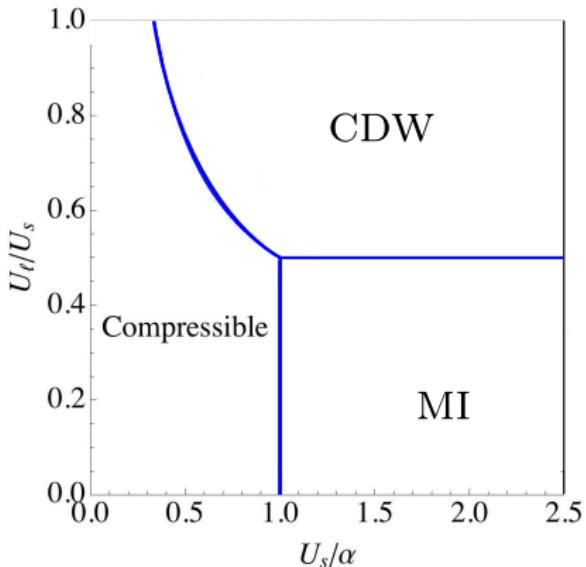
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- \exists compressible phase
 - no distinction: SS & SF
 - MI - SF: off by factor 2
- MI - CDW: accurate
- similar to Gutzwiller and QMC
Batrouni et al '17

Discrete WKB method

Procedure

- effective Hamiltonian

$$\langle \psi | H | \psi \rangle = \sum_Q \epsilon_Q \psi_Q^* \psi_Q + \gamma_Q^+ \psi_{Q+2}^* \psi_Q + \gamma_Q^- \psi_{Q-2}^* \psi_Q$$

- define momenta:

$$\cos p(Q) = \frac{E - \epsilon_Q}{2\gamma_Q}$$

- with $\gamma_Q := \frac{\gamma_Q^+ + \gamma_Q^-}{2}$
- effective tight-binding model
- **classically allowed regions:**

$$p(Q) \in \mathbb{R}$$

$$\Rightarrow \epsilon_Q + 2\gamma_Q \leq E \leq \epsilon_Q - 2\gamma_Q$$

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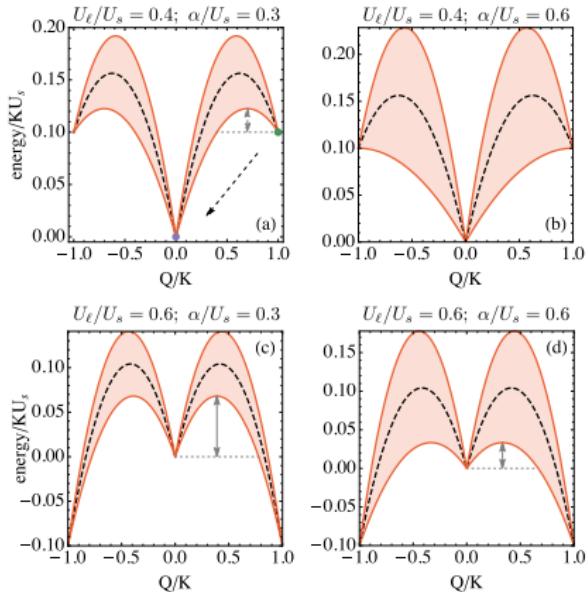
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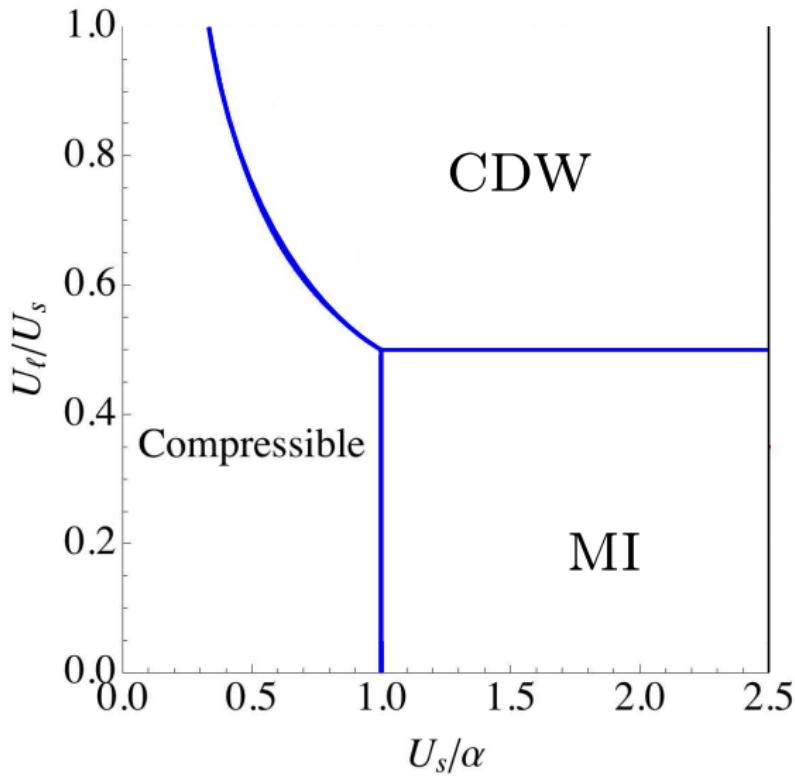
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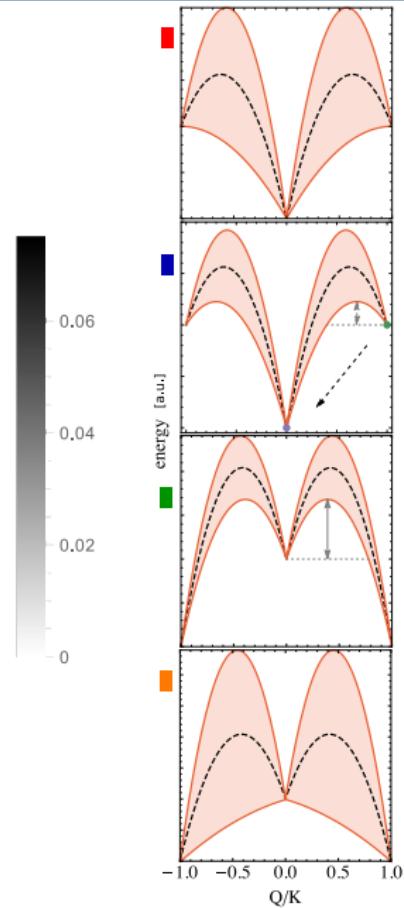
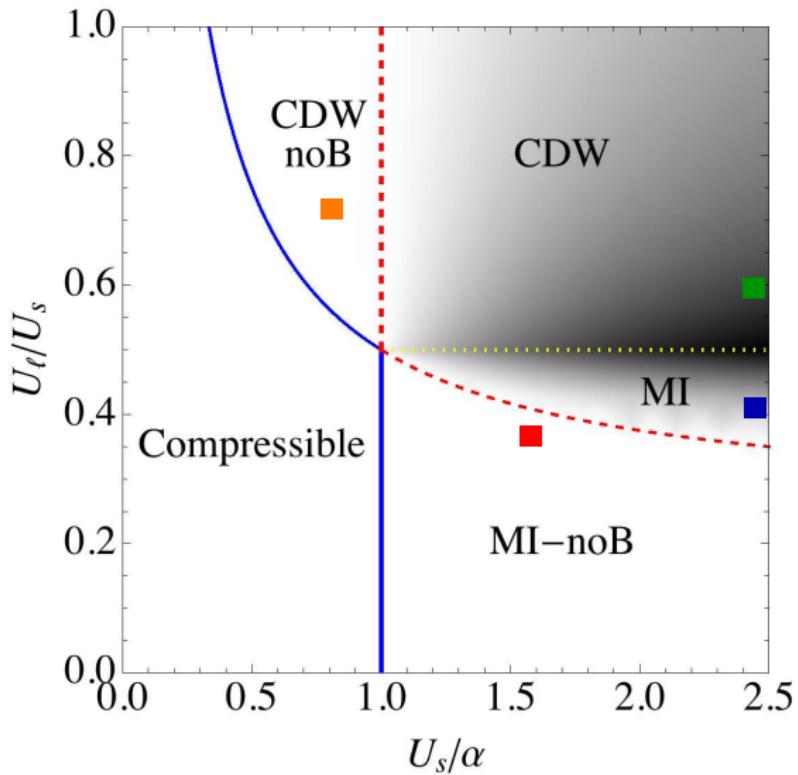
- hopping: simplifies transition

$$\Delta\mathcal{E} = \frac{K/8}{2U_\ell - \alpha} \begin{cases} (U_s - \alpha)^2, & \text{CDW} \\ (\alpha + U_s - 4U_\ell)^2, & \text{MI} \end{cases}$$

Augmented phase diagram



Augmented phase diagram



Conclusion

- role of quantum fluctuations in phase reconfig. of 1st order transition
- long-range Bose-Hubbard model: MI - CDW transition
 - ultra-cold Ru atoms in optical lattice + cavity
 - full control on many-body properties
- exact Landau theory at zero hopping: hysteresis experiments
 - MI phase: meta stable + protected by barrier
 - explain: asymmetric hysteresis + plateaus
- variational description: render problem tractable
 - truncated Hilbert space
 - neglect lattice structure (mean-field like)
 - generate states $|Q\rangle$ by ∞ range hopping
 - numerical study: $\Theta, \Delta E, S_{vN}, \chi \Rightarrow$ phase diagram
- discrete WKB
 - construct phase diagram analytically
 - augment phase diagram by energy barrier
 - observe: tunneling lowers energy barrier

Outlook

- study dynamics in reduced Hilbert space: $\mathcal{L} = \langle \psi | (i\partial_t - H) | \psi \rangle$
 - different quench protocols
 - hysteresis loops
- construct a mean-field Hamiltonian
 - exact descriptions possible ?
 - ∞ range interaction \Rightarrow mean-fiel reliable guide ? CDW - MI
- extend or modify variational $|Q\rangle$ states
 - \exists extension that covers more details of the phase diagram ?
 - distinguish between SS and SF
- describe model as open quantum system: Lindblad, Q-Langevin ?
 - cavity loss
 - incoherent scattering for long-range interaction