

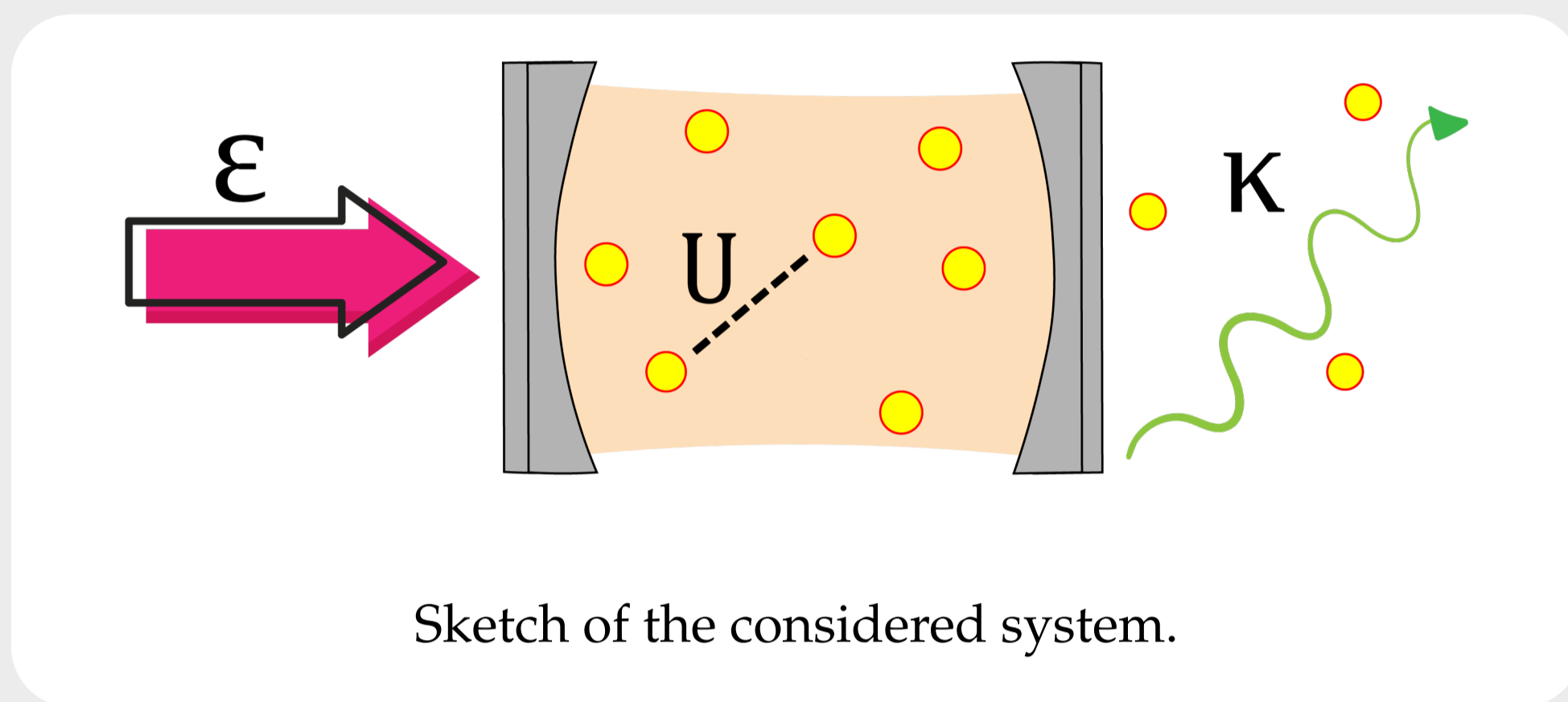
Entropic dynamics of a discontinuous dissipative phase transition

Bruno O. Goes, Institute of Physics, University of São Paulo
Gabriel T. Landi, Institute of Physics, University of São Paulo



Optical bistability model

We are going to present our results for the optical bistability model, which is known to present a discontinuous phase transition in the thermodynamic limit.



The hamiltonian that describes a single-mode field inside a cavity which contains a nonlinear dispersive medium, in the rotating frame is¹

$$H = \Delta_{cp} a^\dagger a + i\mathcal{E}(a^\dagger - a) + \frac{U}{2} a^\dagger a^\dagger a a$$

- a^\dagger (a) creates (annihilates) an excitation inside the cavity;
- $\Delta_{cp} = \omega_c - \omega_p$ is the detuning;
- \mathcal{E} is the amplitude of the drive;
- U is the photon-photon interaction strength.

The one photon losses can be described within Born-Markov approximation, so the density matrix dynamics of the system can be described in terms of a Lindblad master equation:

$$\partial_t \rho = -i[H, \rho] + 2\kappa \left(a \rho a^\dagger - \frac{1}{2} \{a^\dagger a, \rho\} \right) \quad (1)$$

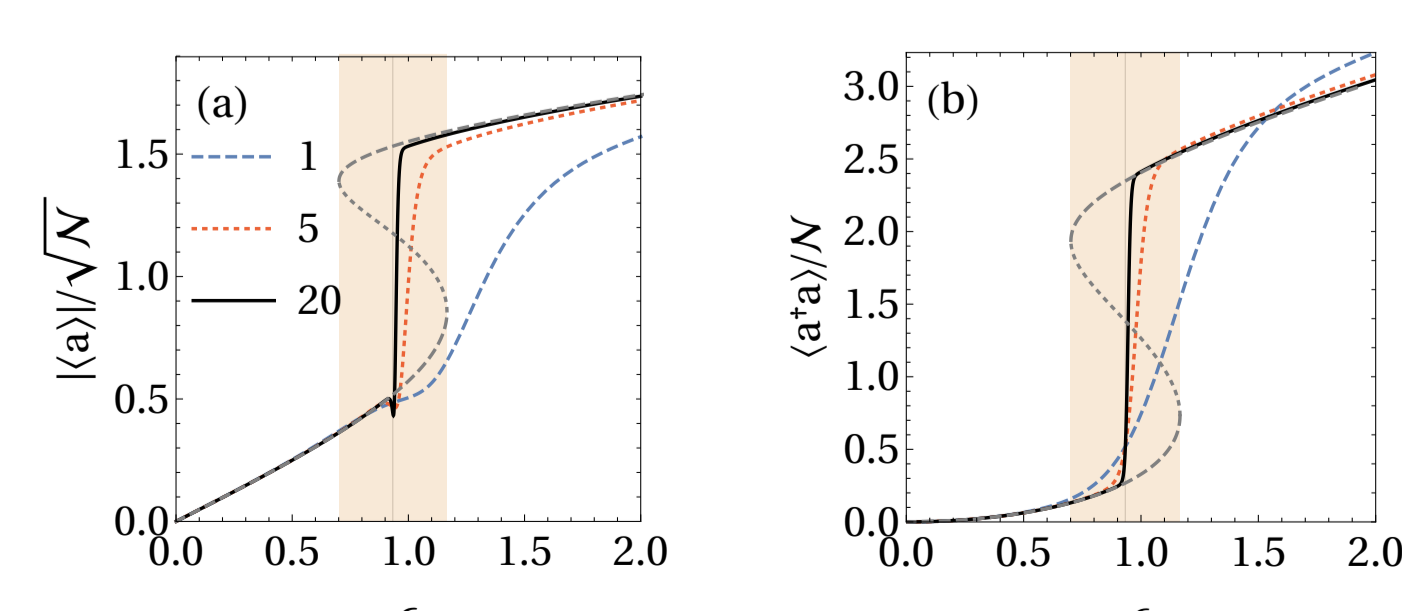
- κ being the dissipation rate.
- Due to the competition between the unitary dynamics and the dissipation a phase transition occurs for a given value ϵ_c .

We introduce the dimensionless parameter N and perform the scaling²

$$U = u/N \quad \mathcal{E} = \sqrt{N}\epsilon \quad \alpha = \sqrt{N}\beta \quad (2)$$

- The thermodynamic limit (TL) occurs when $N \rightarrow \infty$, where $\mathcal{E}^2 U = \text{constant}$ while the number of excitations $n = |\alpha|^2 \propto N$ grows up;
- In the TL quantum fluctuations became negligible and the system behaves in a semiclassical fashion, which justifies a mean field approximation (MFA), which gives us:

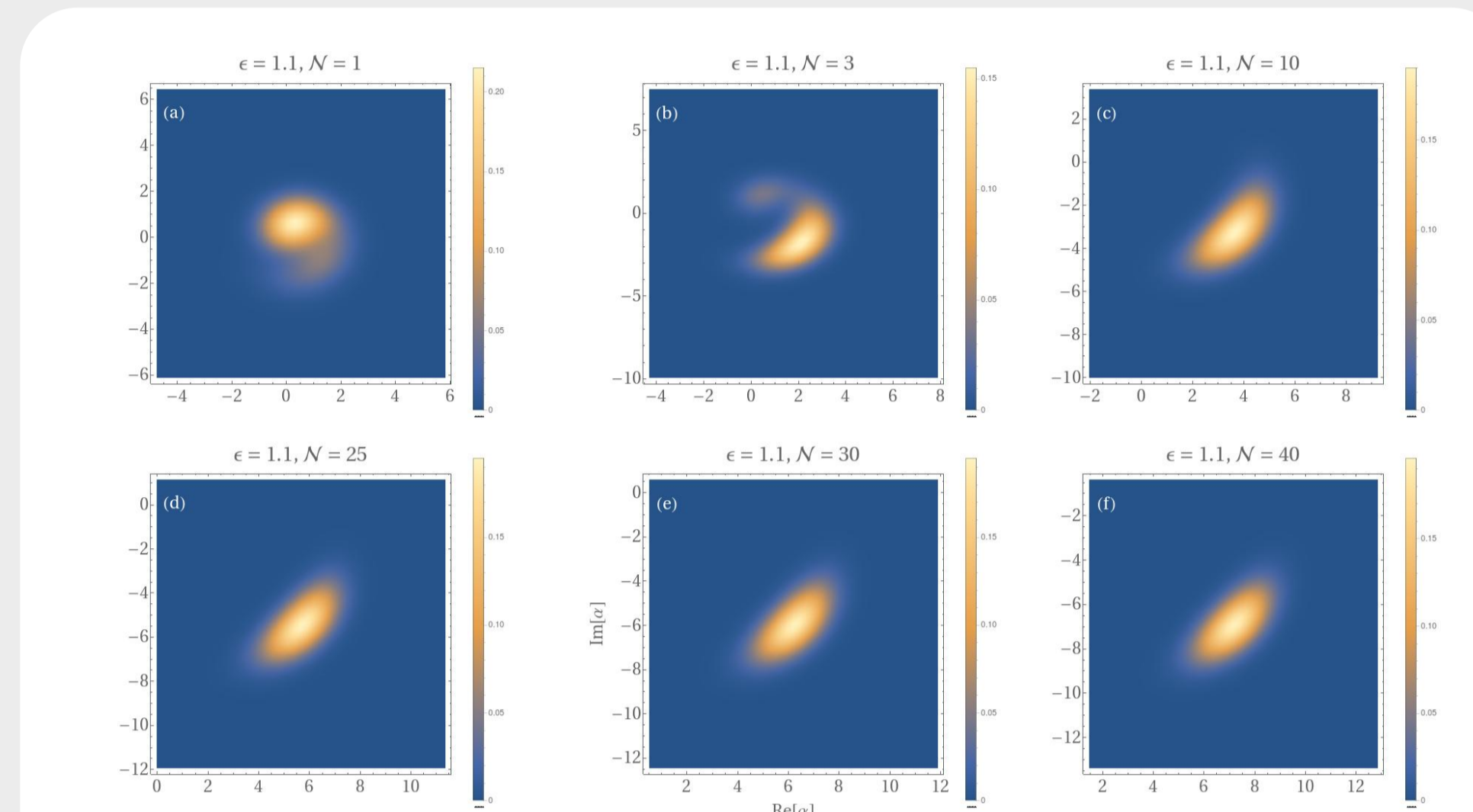
$$\partial_t \alpha = (\kappa - i\Delta_{cp} - iU|\alpha|^2)\alpha + \mathcal{E}$$



$\langle a \rangle / \sqrt{N}$ and $\langle a^\dagger a \rangle / N$, parameters were fixed at $\kappa = 1/2$, $\Delta = -2$ and $U = 1$.

Q-function

- In the quantum phase space the Q-function $Q(\alpha, \alpha^*) = \langle \alpha | \rho | \alpha \rangle / \pi$ tends to a gaussian.



Husimi Q-function contour plot: $Q = \langle \alpha | \rho | \alpha \rangle / \pi$ for $\epsilon = 1.1$ the values of N are (a) $N = 1$ (b) $N = 3$ (c) $N = 10$ (d) $N = 25$ (e) $N = 30$ (f) $N = 40$. Other parameters were fixed at $\kappa = 1/2$, $\Delta = -2$ and $U = 1$.

The Fokker-Planck equation and Wehrl entropy

- We can map eqn. 1 into a FP eqn.:^{3,4}

$$\partial_t Q(\alpha, \alpha^*) = \mathcal{U}(Q) + \mathcal{D}(Q)$$

- And defining the entropy:

$$S_Q = - \int d^2\alpha Q(\alpha, \alpha^*) \ln(Q(\alpha, \alpha^*))$$

$$\frac{dS_Q}{dt} = - \int d^2\alpha \mathcal{U}(Q) \ln(Q) - \int d^2\alpha \mathcal{D}(Q) \ln(Q) = \Psi + (\Pi - \Phi)$$

$$\Psi = \Psi_{\Delta_{cp}} + \Psi_{\mathcal{E}} + \Psi_U \rightarrow \Psi = \Psi_U$$

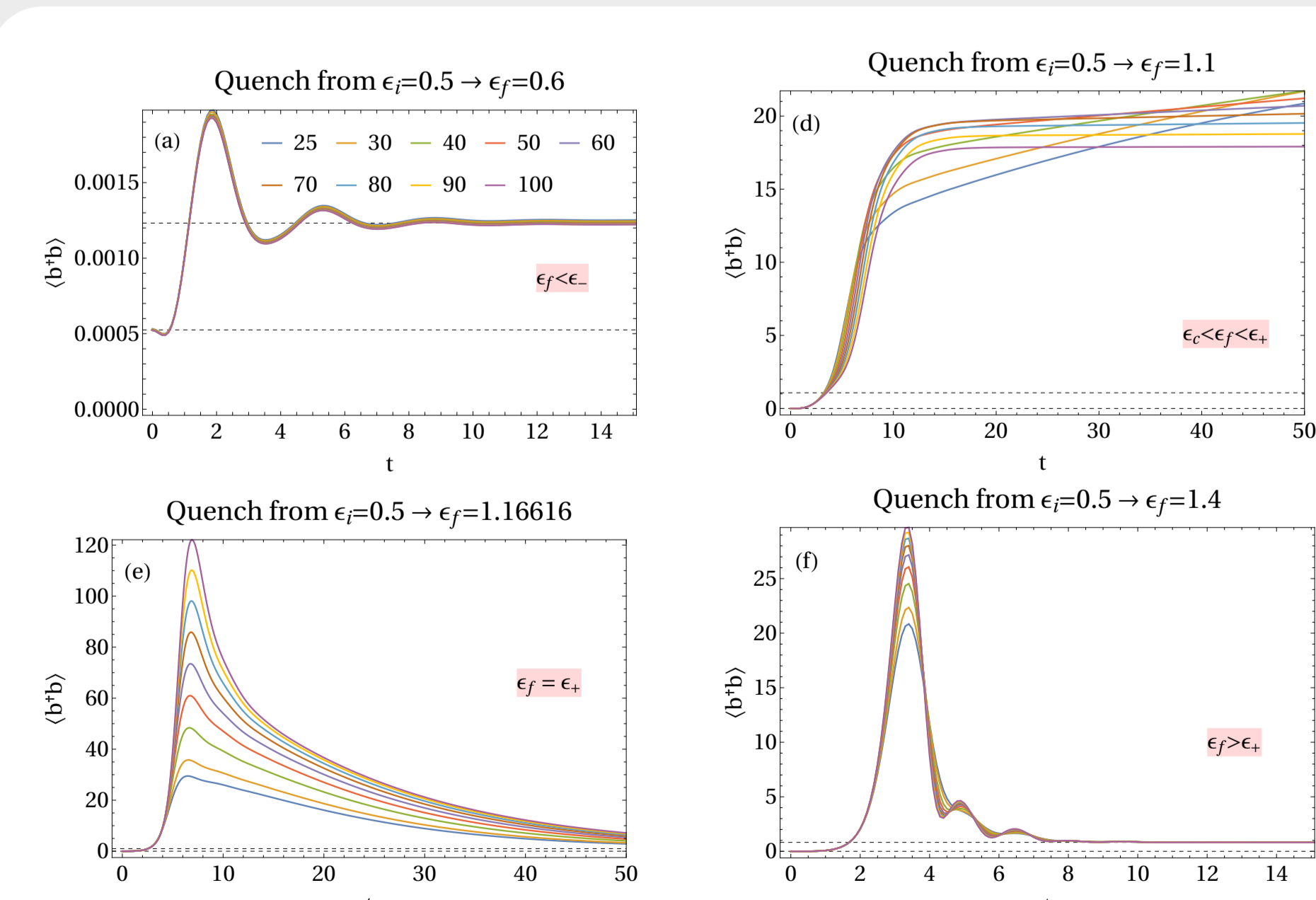
$$\Pi = \frac{2}{\kappa} \int d^2\alpha \frac{|J(Q)|^2}{Q}$$

$$J(Q) = \kappa(\alpha Q + \partial_{\alpha^*} Q)$$

$$\Phi = 2\kappa \langle a^\dagger a \rangle = 2\kappa(|\alpha|^2 + \langle b^\dagger b \rangle) = \Phi_M + \Phi_Q$$

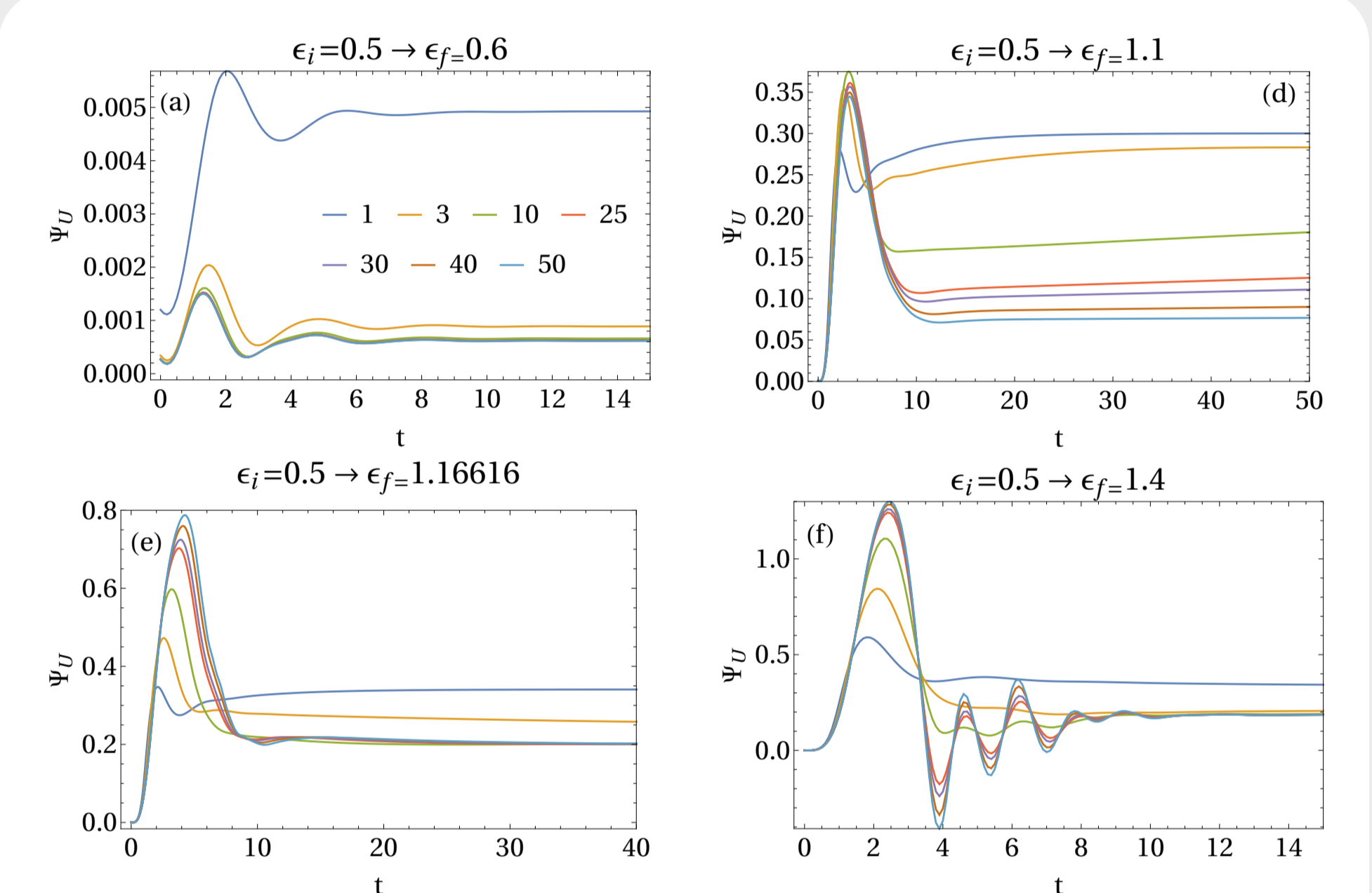
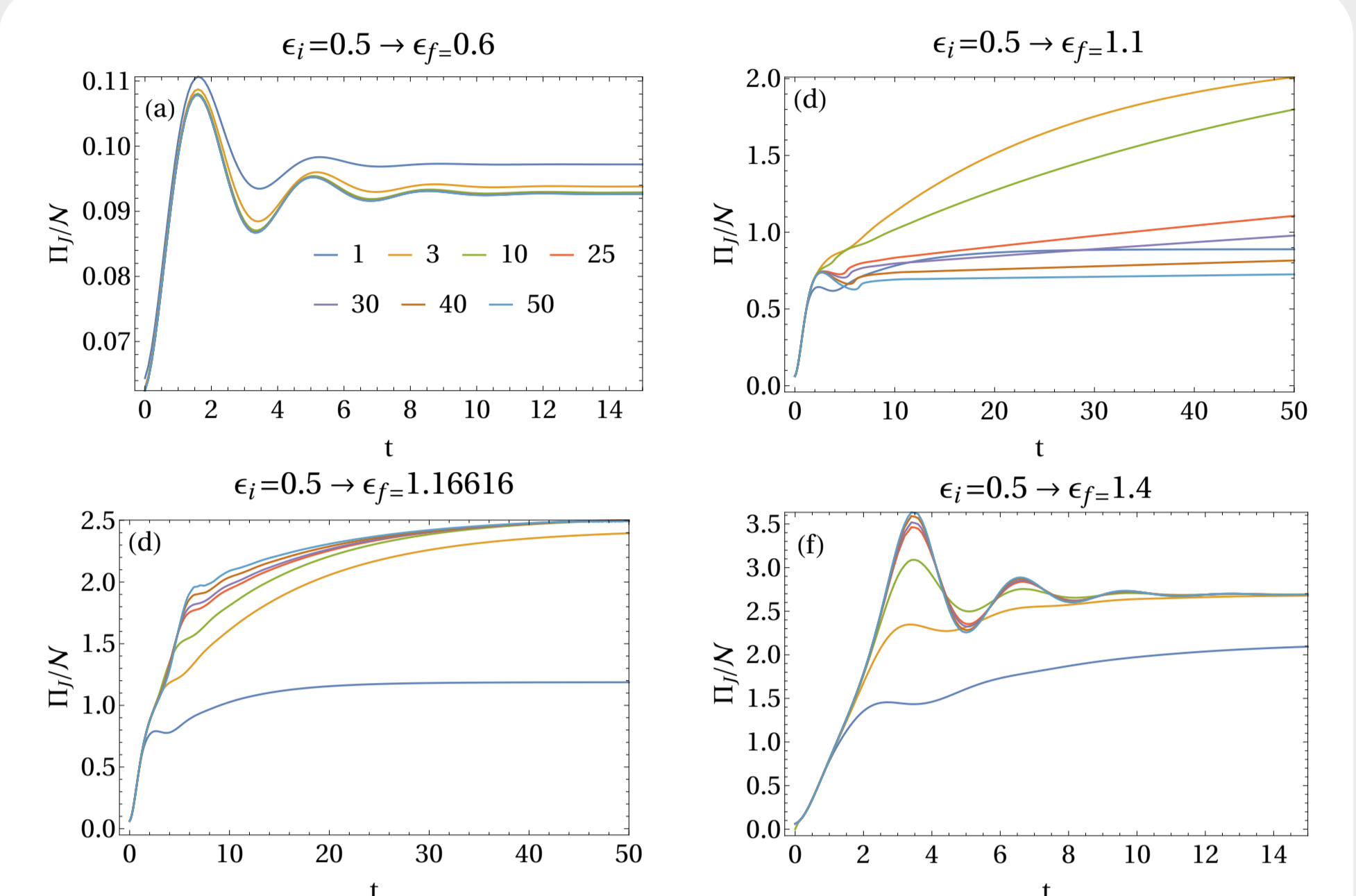
$$\Psi = \frac{iU}{2} \langle \partial_{\alpha^*} ((\alpha^*)^2 \partial_{\alpha^*} (\ln(Q))) - h.c. \rangle$$

- Ψ and Φ_Q does not depend on N ;



Quantum entropy flux $\Phi_Q = 2\kappa \langle b^\dagger b \rangle_t$ for a quench dynamics scenario as a function of time t : The curves are for nine different values of N , as shown in image (a). The dashed black horizontal lines corresponds to NESS value of $\langle b^\dagger b \rangle$ for each pair of (ϵ_i, ϵ_f) , computed from the exact solution of Ref.¹ The initial ϵ_i is fixed at $\epsilon_i = 0.5$. (a) $\epsilon_f = 0.6$ (d) $\epsilon_f = 1.1$ (e) $\epsilon_f = \epsilon_p$ (f) $\epsilon_f = 1.4$. Other parameters were fixed at $\kappa = 1/2$, $\Delta = -2$ and $U = 1$.

Entropic dynamics



Π_U/N and Ψ_U entropic dynamics: The curves are for seven different values of N , as shown in images (a). The initial ϵ_i is fixed $\epsilon_i = 0.5$. (a) $\epsilon_f = 0.6$ (d) $\epsilon_f = 1.1$ (e) $\epsilon_f = \epsilon_p$ (f) $\epsilon_f = 1.4$. Other parameters were fixed at $\kappa = 1/2$, $\Delta = -2$ and $U = 1$.

Conclusions

- We have developed a theoretical framework based on the Husimi Q-function to study the entropic dynamics of a dissipative phase transition, which have 3 contributions;
- We have simulated the dynamics of the system in a quench scenario and we can note that even quantities that does not depend on N , in principle, shows a scaling behaviour when crossing the ϵ_c in the transient dynamics.

Acknowledgements

The authors acknowledge the financial support from CNPq.

References

- 1 P. Drummond and D. Walls, Quantum theory of optical bistability. i. nonlinear polarisability model, Journal of Physics A: Mathematical and General **13**, 725 (1980).
- 2 W. Casteels, R. Fazio, and C. Ciuti, Critical dynamical properties of a first-order dissipative phase transition, Phys. Rev. A **95**, 012128 (2017).
- 3 J. P. Santos, G. T. Landi, and M. Paternostro, Wigner entropy production rate, Phys. Rev. Lett. **118**, 220601 (2017).
- 4 C. Gardiner and P. Zoller, *Quantum noise: a handbook of Markovian and non-Markovian quantum stochastic methods with applications to quantum optics*, volume 56 (Springer Science & Business Media 2004).