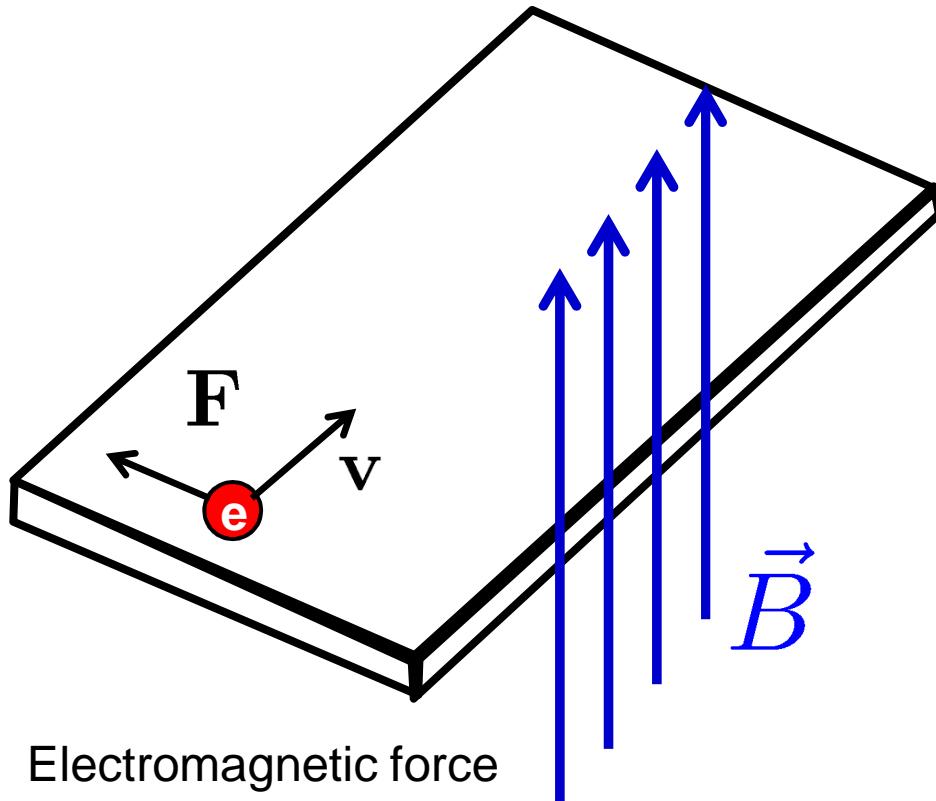


Velocity and Berry curvature in the QHE



Electromagnetic force

$$\mathbf{F} = (-e)\mathbf{E} + (-e)\mathbf{v} \times \mathbf{B}$$

Electric potential and Vector potential.

$$\left\{ \begin{array}{l} \mathbf{B} = \nabla_{\mathbf{r}} \times \mathbf{A}(\mathbf{r}) \\ \mathbf{E} = -\nabla_{\mathbf{r}} V(\mathbf{r}) - \frac{\partial}{\partial t} \mathbf{A}(\mathbf{r}, t) \end{array} \right.$$

In the absence of other charges:

$$V(\mathbf{r}) = 0 \Rightarrow \mathbf{E} = -\frac{\partial}{\partial t} \mathbf{A}(\mathbf{r}, t)$$

Hamiltonian:

$$H = \frac{(\mathbf{p} + e\mathbf{A}(t))^2}{2m^*} = \frac{\hbar^2 (\mathbf{k}(t))^2}{2m^*}$$

Velocity: $m^* \mathbf{v}(t) = \mathbf{p} + e\mathbf{A}(t) = \hbar \mathbf{k}(t)$

Thus: $\mathbf{v}(\mathbf{k}) = \frac{1}{\hbar} \nabla_{\mathbf{k}} H(\mathbf{k})$

Tarefa 18: identity for the velocity

Using:

$$\left\{ \begin{array}{l} \mathbf{v}(\mathbf{k}) = \frac{1}{\hbar} \nabla_{\mathbf{k}} H(\mathbf{k}) \\ H|n, \mathbf{k}(t)\rangle = E_n[\mathbf{k}(t)]|n, \mathbf{k}(t)\rangle \\ i\hbar \frac{d}{dt}|n, \mathbf{k}(t)\rangle = i\hbar \frac{d\mathbf{k}(t)}{dt} \cdot \nabla_{\mathbf{k}}|n, \mathbf{k}(t)\rangle \\ H|\Psi(t)\rangle = i\hbar \frac{d}{dt}|\Psi(t)\rangle \\ \nabla_{\mathbf{k}}(H|n, \mathbf{k}\rangle) = (\nabla_{\mathbf{k}}H)|n, \mathbf{k}\rangle + H(\nabla_{\mathbf{k}}|n, \mathbf{k}\rangle) \end{array} \right.$$

Show that:

$$\hbar \mathbf{v}|n, \mathbf{k}(t)\rangle = \nabla_{\mathbf{k}}(E_n[\mathbf{k}]|n, \mathbf{k}\rangle) - i\hbar \frac{d\mathbf{k}(t)}{dt} \cdot \nabla_{\mathbf{k}}(\nabla_{\mathbf{k}}|n, \mathbf{k}\rangle)$$

Velocity and Berry curvature in the QHE

From the previous result, it follows(*):

$$\mathbf{v}_n(\mathbf{k}) = \langle n, \mathbf{k}(t) | \mathbf{v} | n, \mathbf{k}(t) \rangle = \frac{1}{\hbar} \nabla_{\mathbf{k}} E[\mathbf{k}] + \frac{d\mathbf{k}(t)}{dt} \times \nabla_{\mathbf{k}} \times \langle n, \mathbf{k} | i \nabla_{\mathbf{k}} | n, \mathbf{k} \rangle$$

Remember the definition of the Berry curvature:

$$\Omega_n(\mathbf{k}) = \nabla_{\mathbf{k}} \times \langle n, \mathbf{k} | i \nabla_{\mathbf{k}} | n, \mathbf{k} \rangle$$

and using:

$$\frac{d\mathbf{k}(t)}{dt} = \frac{e}{\hbar} \frac{\partial \mathbf{A}}{\partial t} = -\frac{e}{\hbar} \mathbf{E}$$

we get

$$\boxed{\mathbf{v}_n(\mathbf{k}) = \frac{1}{\hbar} \nabla_{\mathbf{k}} E_n[\mathbf{k}] - \frac{e}{\hbar} \mathbf{E} \times \Omega_n(\mathbf{k})}$$

(*) Should be on Lista 5!

Hall Conductance and Chern number

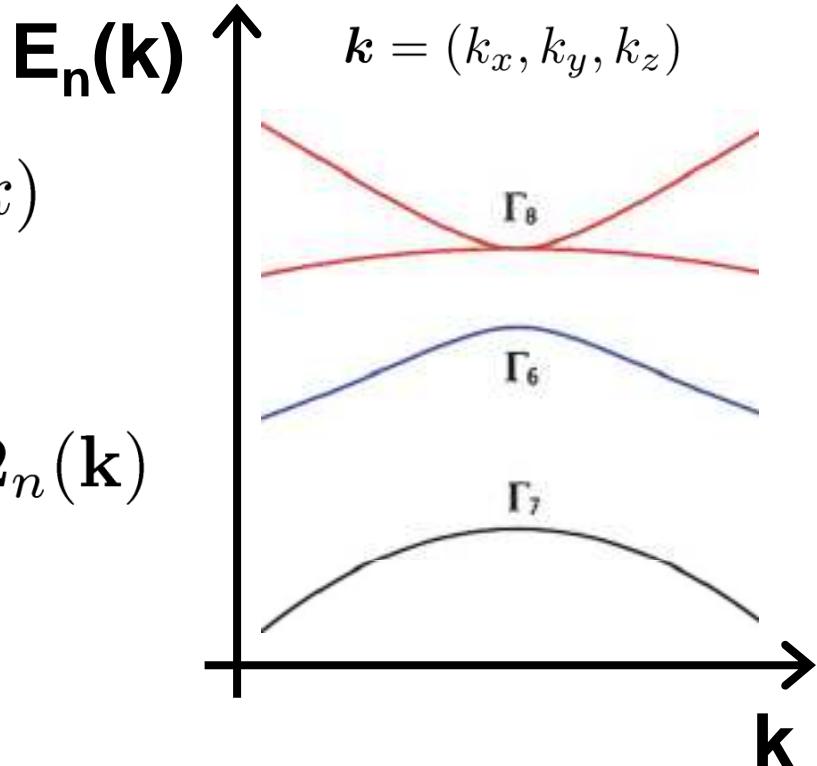
Current density(*): $\mathbf{J} = -e \sum_n \int \frac{d\mathbf{k}}{(2\pi)^2} \mathbf{v}_n(\mathbf{k}) f(k)$

Using: $\mathbf{v}_n(\mathbf{k}) = \frac{1}{\hbar} \nabla_{\mathbf{k}} E_n[\mathbf{k}] - \frac{e}{\hbar} \mathbf{E} \times \boldsymbol{\Omega}_n(\mathbf{k})$

we might calculate the conductance: $\mathbf{J} = \boldsymbol{\sigma} \cdot \mathbf{E}$

If we have a gap and N filled levels $\sum_{n \in \text{filled}} \int \frac{d\mathbf{k}}{(2\pi)^2} \nabla_{\mathbf{k}} E_n[\mathbf{k}] f(k) = 0$

(*) Quantum version of the usual: $\mathbf{J} = (-e)n\langle \mathbf{v} \rangle$



Hall Conductance and Chern number

The conductance can then be calculated: $\mathbf{J} = \boldsymbol{\sigma} \cdot \mathbf{E}$

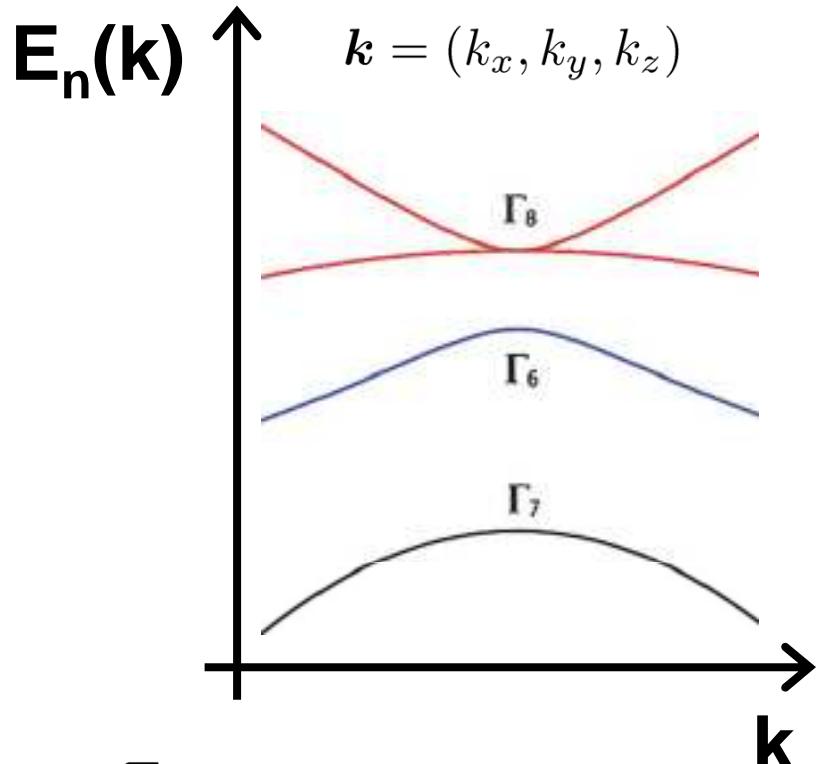
$$\sigma_{xy} = \frac{e^2}{h} \frac{1}{2\pi} \sum_{n \in \text{filled}} \int_{\text{BZ}} \boldsymbol{\Omega}_n(\mathbf{k}) \cdot d\mathbf{k}$$

The integral will be carried out in the 1st BZ, which is a torus for the Berry curvature:

$$\boldsymbol{\Omega}_n(k_x, k_y) = \boldsymbol{\Omega}_n(k_x + \frac{\pi}{a}, k_y) = \boldsymbol{\Omega}_n(k_x, k_y + \frac{\pi}{a})$$

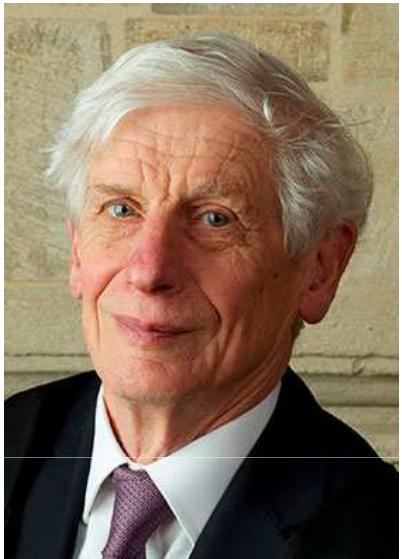
Thus the integral will be 2π (Chern number) and the sum will give the number of filled bands ν :

$$\boxed{\sigma_{xy} = \frac{e^2}{h} \frac{1}{2\pi} 2\pi \nu = \frac{e^2}{h} \nu}$$



TKNN invariant: 1982

The Hall conductivity is proportional to a Chern number (Berry-phase-like)



David Thouless

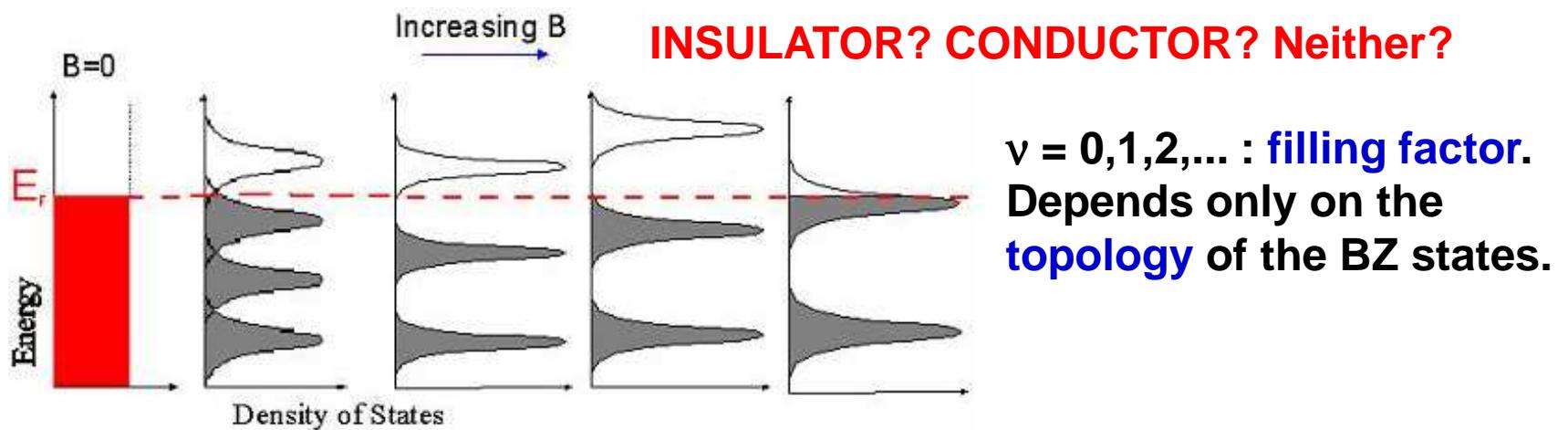


2016

$$\sigma_{xy} = \frac{e^2}{h} \sum_{n < N_F} \frac{1}{2\pi} \iint_{BZ} \Omega_n(\mathbf{k}) \cdot d\mathbf{k} \equiv \nu \frac{e^2}{h}$$

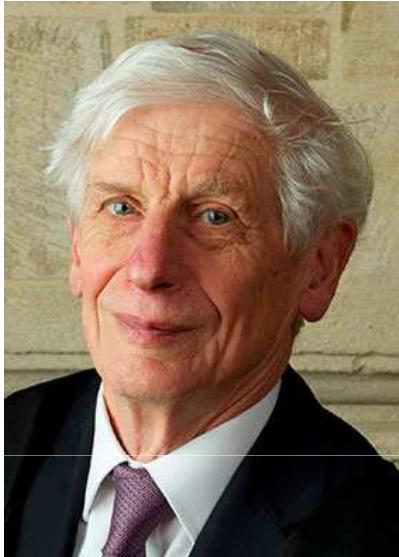
Thouless, Kohmoto, Nightingale, den Nijs, *Phys. Rev. Lett.* 49, 405 (1982)

- System is periodic (BZ is a torus in k-space)
- There is an uniform magnetic field in the system.
- Fermi energy lies in a gap with N_F filled bands.



TKNN invariant: 1982

The Hall conductivity is proportional to a Chern number (Berry-phase-like)



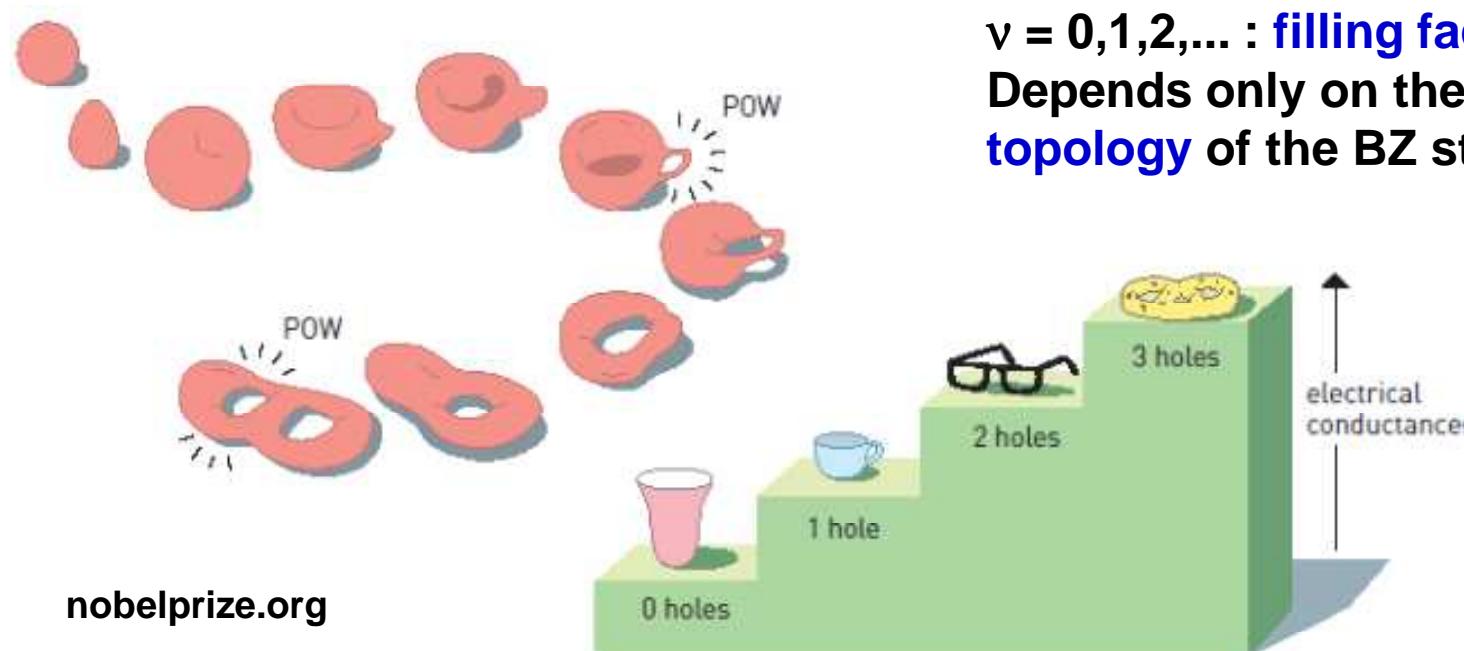
David Thouless



2016

$$\sigma_{xy} = \frac{e^2}{h} \sum_{n < N_F} \frac{1}{2\pi} \iint_{BZ} \Omega_n(\mathbf{k}) \cdot d\mathbf{k} \equiv \nu \frac{e^2}{h}$$

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