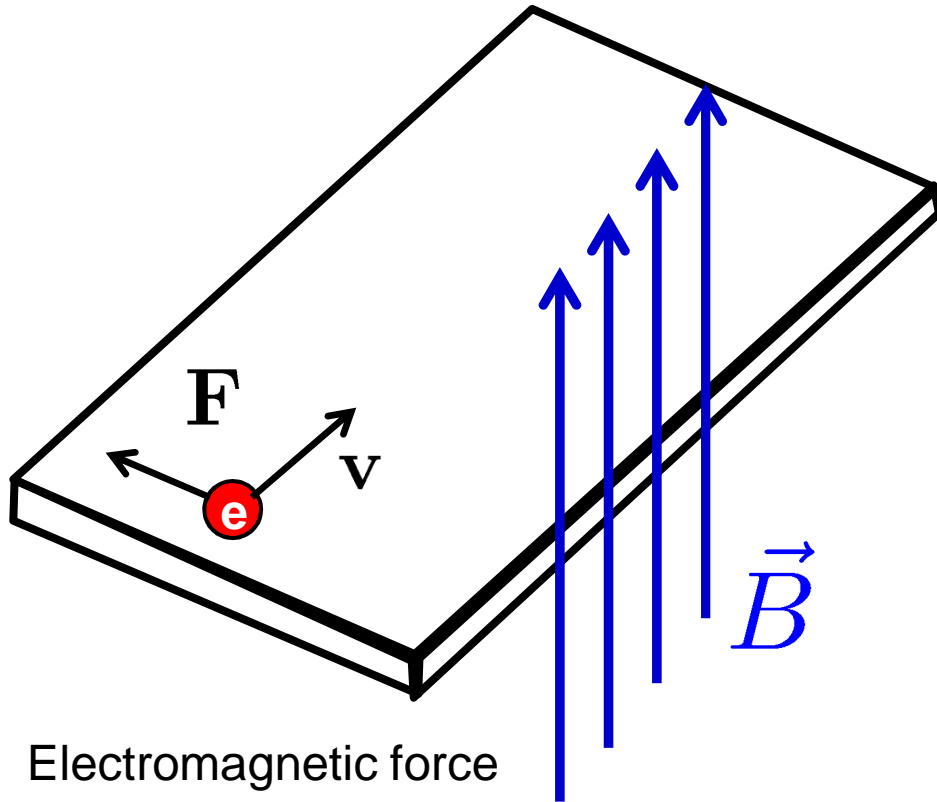


# Velocity and Berry curvature in the QHE



Electromagnetic force

$$\mathbf{F} = (-e)\mathbf{E} + (-e)\mathbf{v} \times \mathbf{B}$$

Electric potential and Vector potential.

$$\begin{cases} \mathbf{B} = \nabla_{\mathbf{r}} \times \mathbf{A}(\mathbf{r}) \\ \mathbf{E} = -\nabla_{\mathbf{r}} V(\mathbf{r}) - \frac{\partial}{\partial t} \mathbf{A}(\mathbf{r}, t) \end{cases}$$

In the absence of other charges:

$$V(\mathbf{r}) = 0 \Rightarrow \mathbf{E} = -\frac{\partial}{\partial t} \mathbf{A}(\mathbf{r}, t)$$

**Hamiltonian:**

$$H = \frac{(\mathbf{p} + e\mathbf{A}(t))^2}{2m^*} \equiv \frac{\hbar^2 (\mathbf{k}(t))^2}{2m^*}$$

**Velocity:**  $m^* \mathbf{v}(t) = \mathbf{p} + e\mathbf{A}(t) = \hbar \mathbf{k}(t)$

**Thus:**  $\mathbf{v}(\mathbf{k}) = \frac{1}{\hbar} \nabla_{\mathbf{k}} H(\mathbf{k})$

# Tarefa 20: identity for the velocity

Using:

$$\left\{ \begin{array}{l} \mathbf{v}(\mathbf{k}) = \frac{1}{\hbar} \nabla_{\mathbf{k}} H(\mathbf{k}) = \frac{1}{\hbar} \left( \frac{\partial H}{\partial k_x} \mathbf{i} + \frac{\partial H}{\partial k_y} \mathbf{j} + \frac{\partial H}{\partial k_z} \mathbf{k} \right) \\ H|n, \mathbf{k}(t)\rangle = E_n[\mathbf{k}(t)]|n, \mathbf{k}(t)\rangle \\ i\hbar \frac{d}{dt} |n, \mathbf{k}(t)\rangle = i\hbar \frac{d\mathbf{k}(t)}{dt} \cdot \nabla_{\mathbf{k}} |n, \mathbf{k}(t)\rangle \\ H|\Psi(t)\rangle = i\hbar \frac{d}{dt} |\Psi(t)\rangle \end{array} \right.$$

Show that:

$$1) \nabla_{\mathbf{k}} (H|n, \mathbf{k}\rangle) = (\nabla_{\mathbf{k}} H)|n, \mathbf{k}(t)\rangle + H(\nabla_{\mathbf{k}} |n, \mathbf{k}\rangle)$$

$$2) H(\nabla_{\mathbf{k}} |n, \mathbf{k}\rangle) = i\hbar \nabla_{\mathbf{k}} \left( \frac{d\mathbf{k}(t)}{dt} \cdot \nabla_{\mathbf{k}} |n, \mathbf{k}\rangle \right)$$

$$3) \hbar \mathbf{v} |n, \mathbf{k}(t)\rangle = \nabla_{\mathbf{k}} (E_n[\mathbf{k}] |n, \mathbf{k}\rangle) - i\hbar \nabla_{\mathbf{k}} \left( \frac{d\mathbf{k}(t)}{dt} \cdot \nabla_{\mathbf{k}} |n, \mathbf{k}\rangle \right)$$

Tip: Do it by components so you don't get confused!!

# Velocity and Berry curvature in the QHE

From the previous result, it follows(\*):

$$\mathbf{v}_n(\mathbf{k}) = \langle n, \mathbf{k}(t) | \mathbf{v} | n, \mathbf{k}(t) \rangle = \frac{1}{\hbar} \nabla_{\mathbf{k}} E[\mathbf{k}] + \frac{d\mathbf{k}(t)}{dt} \times \nabla_{\mathbf{k}} \times \langle n, \mathbf{k} | i \nabla_{\mathbf{k}} | n, \mathbf{k} \rangle$$

Remember the definition of the Berry curvature:

$$\boldsymbol{\Omega}_n(\mathbf{k}) = \nabla_{\mathbf{k}} \times \langle n, \mathbf{k} | i \nabla_{\mathbf{k}} | n, \mathbf{k} \rangle$$

and using: 
$$\frac{d\mathbf{k}(t)}{dt} = \frac{e}{\hbar} \frac{\partial \mathbf{A}}{\partial t} = -\frac{e}{\hbar} \mathbf{E}$$

we get 
$$\mathbf{v}_n(\mathbf{k}) = \frac{1}{\hbar} \nabla_{\mathbf{k}} E_n[\mathbf{k}] - \frac{e}{\hbar} \mathbf{E} \times \boldsymbol{\Omega}_n(\mathbf{k})$$

(\*) Should be on Lista 5!

# Hall Conductance and Chern number

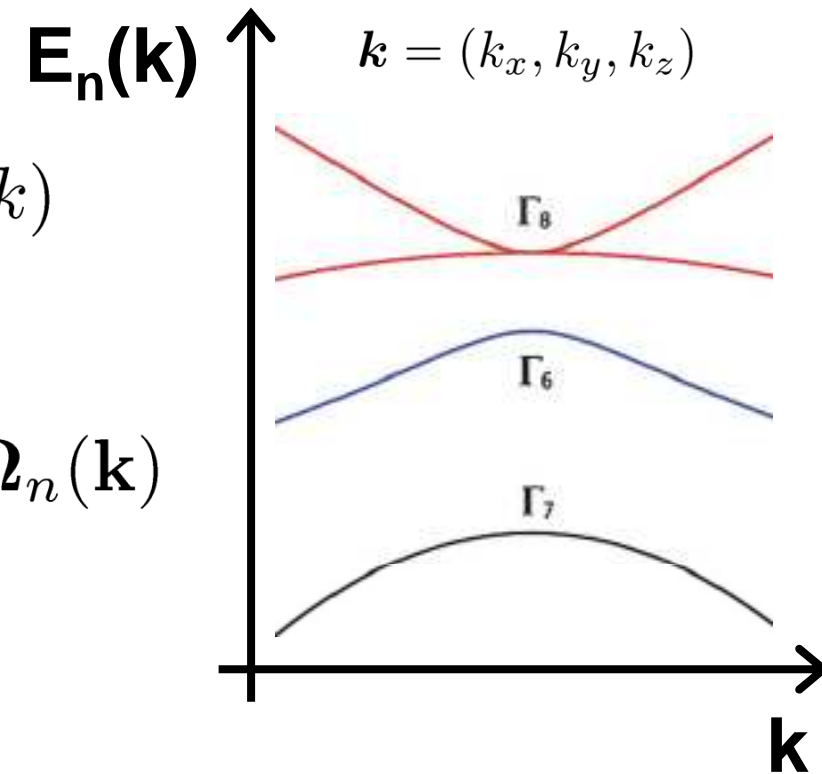
Current density(\*): 
$$\mathbf{J} = -e \sum_n \int \frac{d\mathbf{k}}{(2\pi)^2} \mathbf{v}_n(\mathbf{k}) f(k)$$

Using: 
$$\mathbf{v}_n(\mathbf{k}) = \frac{1}{\hbar} \nabla_{\mathbf{k}} E_n[\mathbf{k}] - \frac{e}{\hbar} \mathbf{E} \times \boldsymbol{\Omega}_n(\mathbf{k})$$

we might calculate the conductance: 
$$\mathbf{J} = \boldsymbol{\sigma} \cdot \mathbf{E}$$

If we have a **gap** and N filled levels 
$$\sum_{n \in \text{filled}} \int \frac{d\mathbf{k}}{(2\pi)^2} \nabla_{\mathbf{k}} E_n[\mathbf{k}] f(k) = 0$$

(\*) Quantum version of the usual: 
$$\mathbf{J} = (-e)n \langle \mathbf{v} \rangle$$



# Hall Conductance and Chern number

The conductance can then be calculated:  $\mathbf{J} = \boldsymbol{\sigma} \cdot \mathbf{E}$

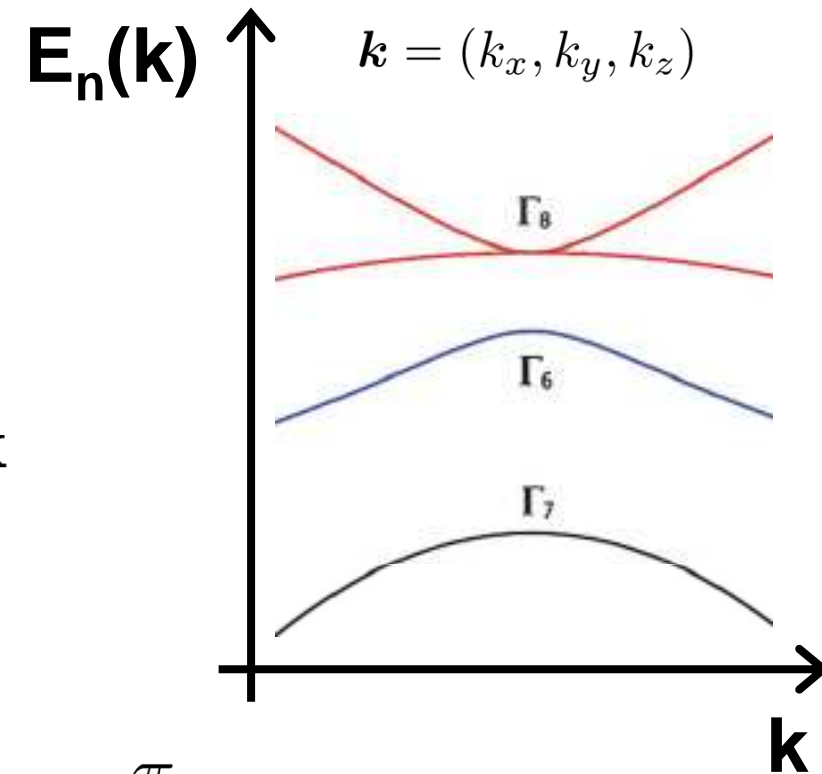
$$\sigma_{xy} = \frac{e^2}{h} \frac{1}{2\pi} \sum_{n \in \text{filled}} \int_{\text{BZ}} \boldsymbol{\Omega}_n(\mathbf{k}) \cdot d\mathbf{k}$$

The integral will be carried out in the 1st BZ, which is a torus for the Berry curvature:

$$\boldsymbol{\Omega}_n(k_x, k_y) = \boldsymbol{\Omega}_n\left(k_x + \frac{\pi}{a}, k_y\right) = \boldsymbol{\Omega}_n\left(k_x, k_y + \frac{\pi}{a}\right)$$

Thus the integral will be  $2\pi$  (Chern number) and the sum will give the number of filled bands  $\nu$  :

$$\sigma_{xy} = \frac{e^2}{h} \frac{1}{2\pi} 2\pi\nu = \frac{e^2}{h} \nu$$



# TKNN invariant: 1982

The Hall conductivity is proportional to a **Chern number** (Berry-phase-like)



$$\sigma_{xy} = \frac{e^2}{h} \sum_{n < N_F} \frac{1}{2\pi} \iint_{\text{BZ}} \Omega_n(\mathbf{k}) \cdot d\mathbf{k} \equiv \nu \frac{e^2}{h}$$

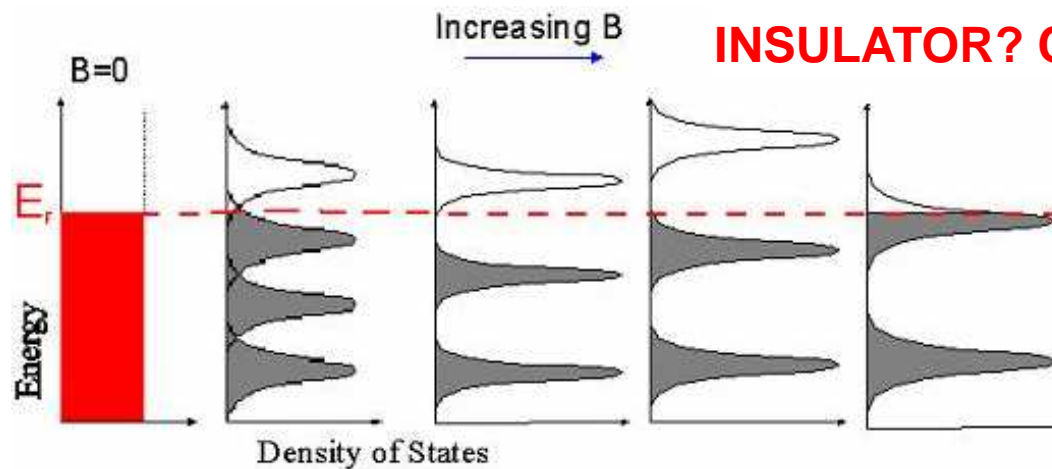
Thouless, Kohmoto, Nightingale, den Nijs, *Phys. Rev. Lett.* 49, 405 (1982)

- System is periodic (BZ is a torus in k-space)
- There is an uniform magnetic field in the system.
- Fermi energy lies in a gap with  $N_F$  filled bands.

David Thouless



2016



$\nu = 0, 1, 2, \dots$  : **filling factor**.  
Depends only on the **topology** of the BZ states.

# TKNN invariant: 1982

The Hall conductivity is proportional to a **Chern number** (Berry-phase-like)



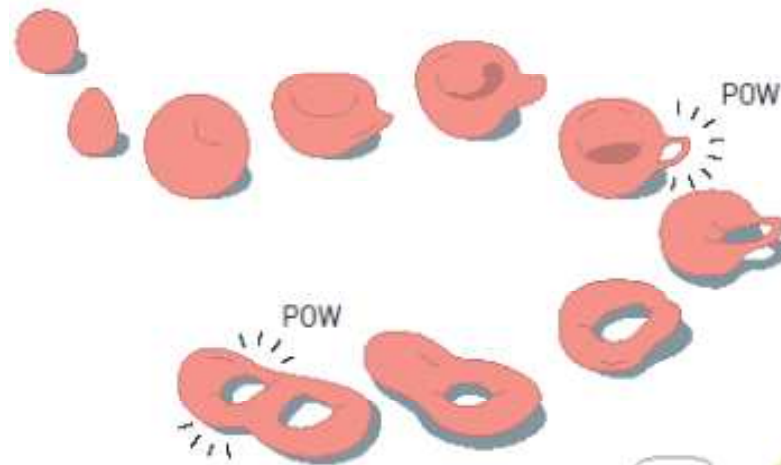
David Thouless



2016

$$\sigma_{xy} = \frac{e^2}{h} \sum_{n < N_F} \frac{1}{2\pi} \iint_{\text{BZ}} \Omega_n(\mathbf{k}) \cdot d\mathbf{k} \equiv \nu \frac{e^2}{h}$$

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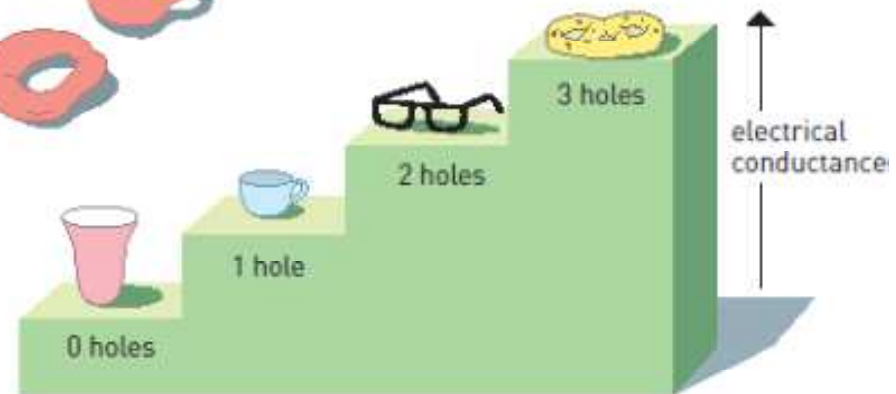


Illustration: ©Johan Jarnestad/The Royal Swedish Academy of Sciences